

Properties of asymmetric nuclear matter

Spectroscopic factors for $N > Z$

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- Symmetry energy
- status of present calculations:
BHF, variational, RMF, χ PT;
Self-consistent Green Function approach
- constraints: Symmetry energy \propto neutron skin
- nuclear matter: spectral functions
- finite nuclei: use of local density approximation (LDA)

Symmetry Energy

proton fraction $x = \rho_p / (\rho_p + \rho_n) = (1 - \alpha) / 2$

$$E(\rho, \alpha) = E(\rho, 0) + S_2(\rho)\alpha^2 + S_4(\rho)\alpha^4 + \dots,$$

$$E(\rho, 0) = E_0 + \frac{K}{18\rho_0^2}(\rho - \rho_0)^2 + \dots,$$

Symmetry energy expanded around saturation density ρ_0

$$S_2(\rho) = \frac{1}{2} \frac{\partial^2 E(\rho, \alpha)}{\partial \alpha^2} \Big|_{\alpha=0} = a_4 + \frac{p_0}{\rho_0^2}(\rho - \rho_0) + \frac{\Delta K}{18\rho_0^2}(\rho - \rho_0)^2 + \dots$$

Note: **Parabolic approximation** is accurate (S_4 small)

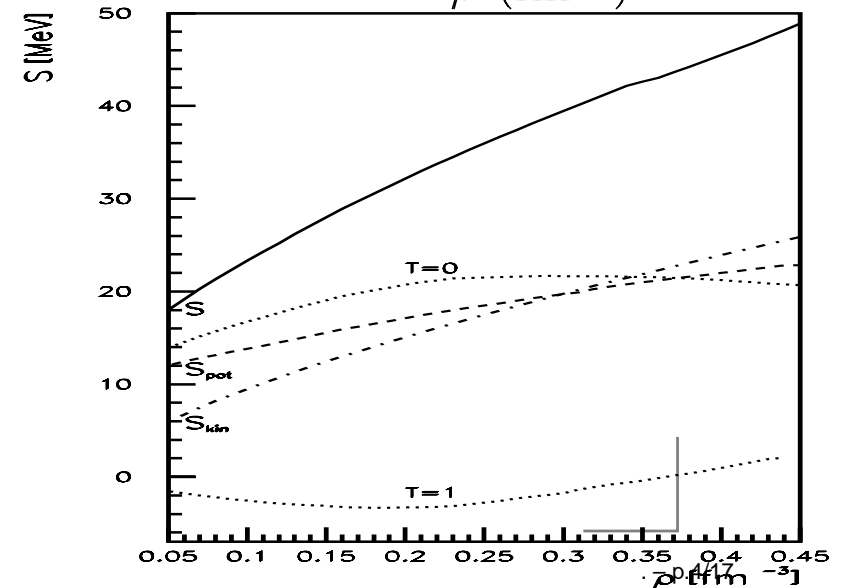
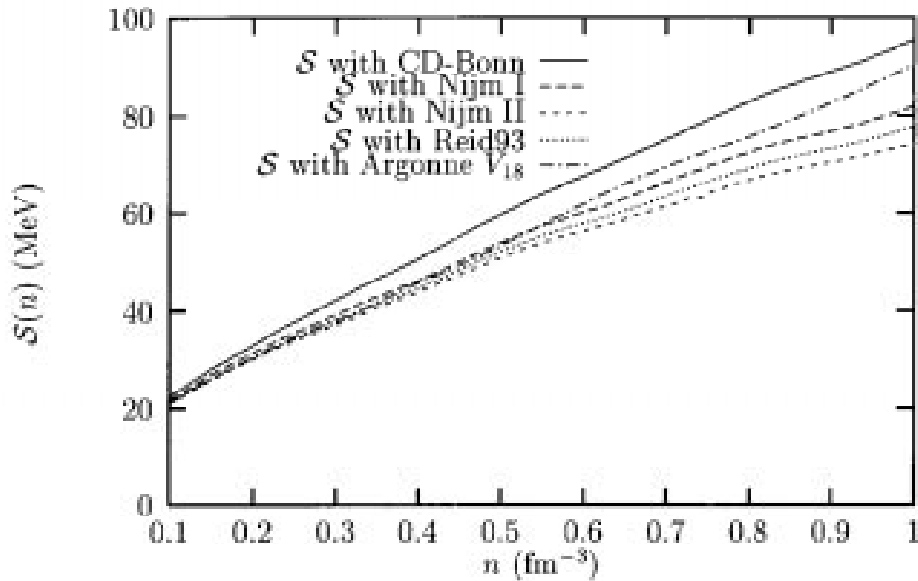
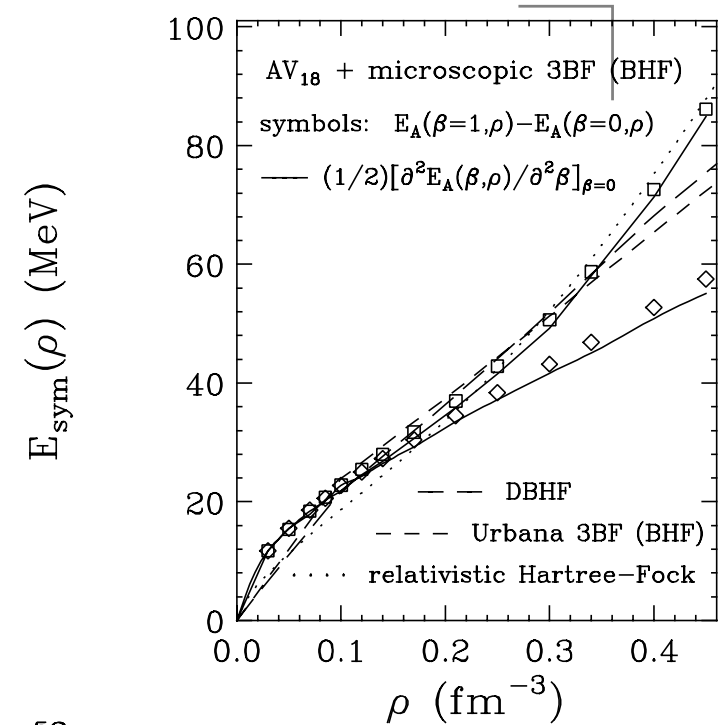
$$a_4 = S_2(\rho_s) = \frac{k_F^2}{6E_F} + S_{\text{pot}}$$

Many-body approaches to EoS

- Realistic NN force as input $\pi + \rho \dots$
Lowest order Brueckner-Hartree-Fock (Dirac-BHF)
Variational approach
Self-consistent Green function approach (New)
EFT
- Effective Lagrangian methods
Skyrme force
Relativistic mean field $\sigma, \omega, \rho \dots$
- Non-conventional
hyperons (input NY and YY interaction)
quark degrees of freedom (QCD)
- Issue: 3-body interactions

(i) Microscopic approaches using NN

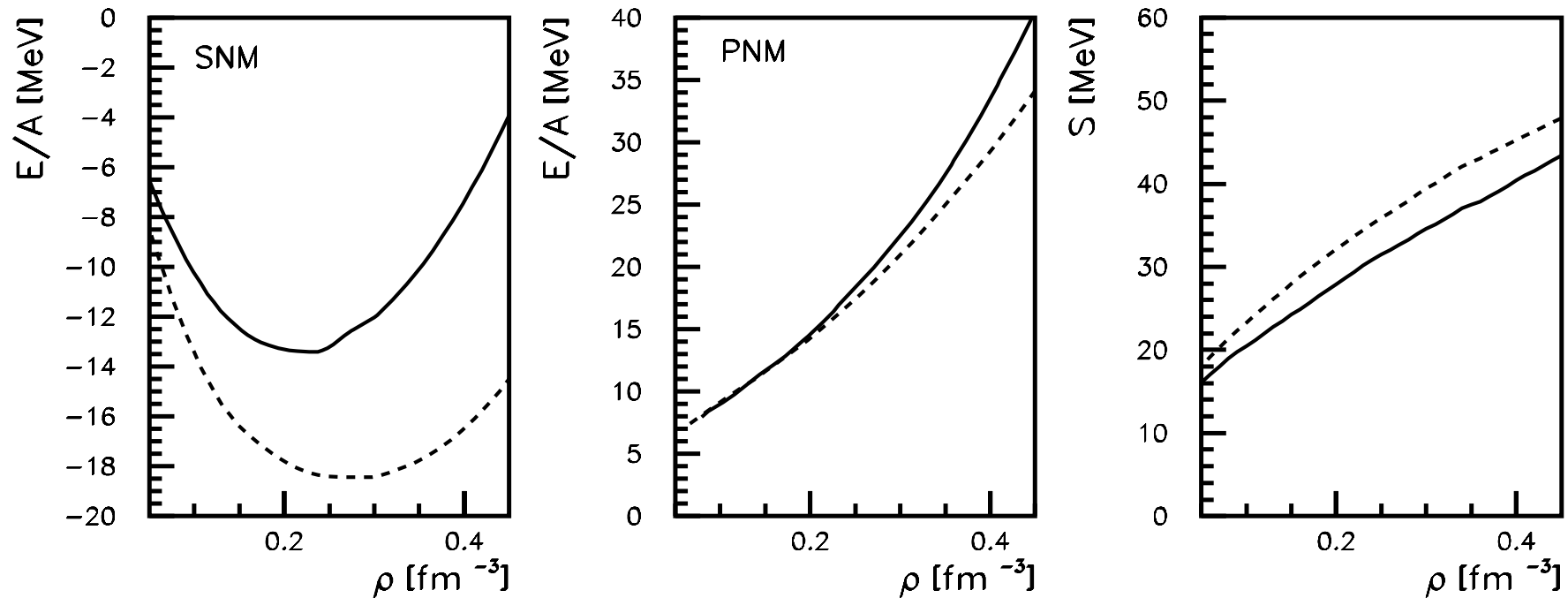
- Lowest Order Brueckner-HF
 - S_{pot} Insensitive to details of NN
 - Tensor force important in S_{pot}
 - effect of 3BF for $\rho > 1.5\rho_s$
- Variational+chain summ. Av14,18
 - Lower S than LOBF
 - 3BF increases stiffness
- Self-Consistent Green Function



SCGF approach

- $T = V + VGGT$ with dressed s.p. propagator
 $G^{-1} = \omega - \frac{k^2}{2m} - \Sigma$; self-energy $i\Sigma = Tr[TG]$
- **Conserving approximation:** Hole and particles treated on equal footing; $\epsilon_F = E_B/A$
- **Binding energy** $\frac{E}{A} = \frac{D}{\rho} \int \frac{d^3k}{(2\pi)^3} \int_{-\infty}^{\epsilon_F} d\omega \left(\frac{k^2}{2m} + \omega \right) S_h(k, \omega)$,
Spectral function:
 $S_h(k, \omega) = ImG(k, \omega) = \langle \psi | a_k^\dagger \delta(\omega - H) a_k | \psi \rangle$
- **Inclusion of high-momentum nucleons:**
SNM: doubling of E_{kin} and E_{pot} , smaller $E(\rho)$;
(increasing with density); **stiffer EoS, improved ρ_s .**
- **PNM:** absence of T=0 tensor force: effect is smaller
Symmetry energy is reduced compared to LOBHF
- **particle-hole correlations not included**

Result SCGF compared to ccBHF



full curve: SCGF; dashed: ccBHF
NN interaction: Reid93, PRC68(2003)

Note: improved saturation density; reduced S_{pot}

(ii) Mean field approaches

- Non-relativistic Skyrme forces
- Relativistic Mean field models: Non-linear extension of Walecka model (**no pions**)
either (i) σ and ω self interactions
or (ii) density dependent couplings
isovector channel: ρ and δ ($I = 1$ **scalar**) exchange

Result: $a_4 \approx \frac{k_F^2}{6(M^{*2} + k_F^2)^{1/2}} + \frac{g_\rho^2}{8m_\rho^2} \rho_s - \frac{g_\delta^2}{8m_\delta^2} \left(\frac{M^*}{\epsilon_F}\right)^2 \rho_s;$

$$p_0 \sim \frac{g_\rho^2}{8m_\rho^2} \rho_s^2$$

Chiral Dynamics

Kaiser et al: Expansion in terms of k_F/m_π up to $O(k_F^5)$
Pion loops with **in-medium** nucleon propagator

$$(\not{p} + M) \left(\frac{i}{p^2 - M^2 + i\epsilon} - 2\pi\delta(p^2 - M^2)\theta(p_0)\theta(k_F - |p|) \right)$$

1π exch Hartree+ Fock with 2+3 medium insertions

iterated one-pion irreducible 2π exchange

- describes SNM with $\Lambda = 0.65$ GeV

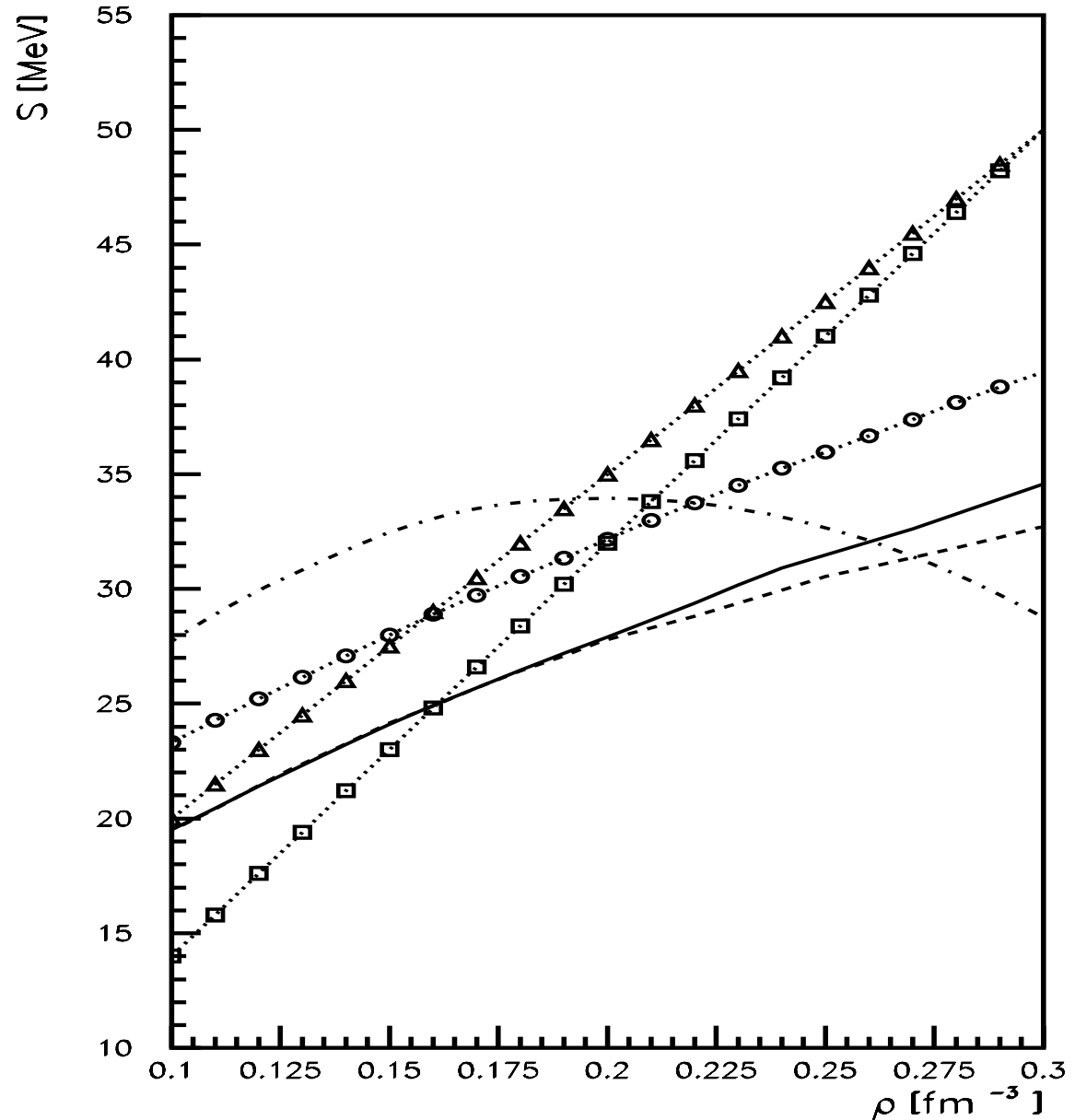
- **chiral limit:**

$$E^{SNM} = \frac{3k_F^2}{10M} - \alpha(\Lambda)\frac{k_F^3}{M^2} + \beta\frac{k_F^4}{M^3}$$

- **SE** levels off at higher ρ

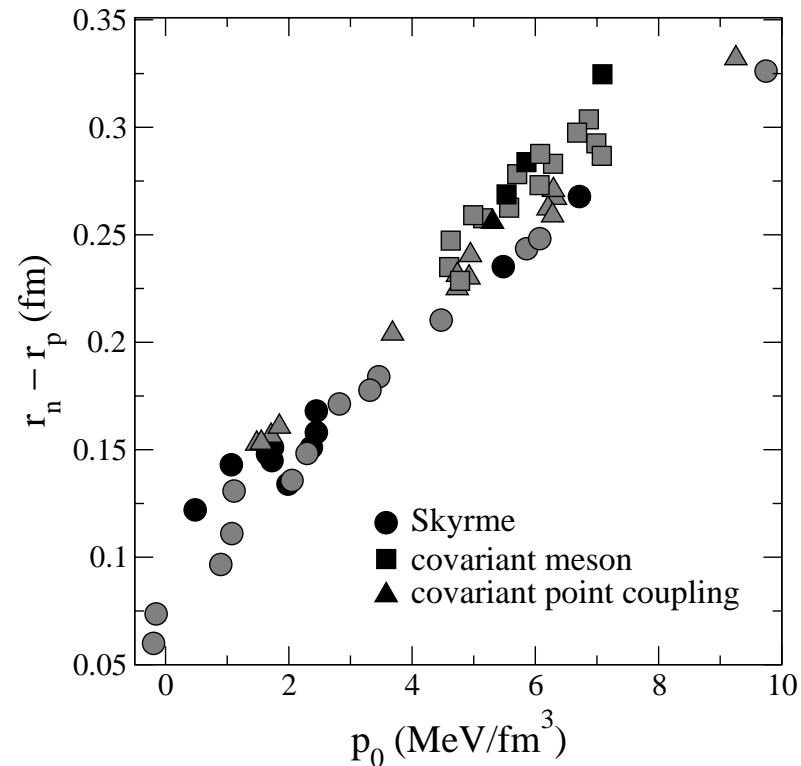
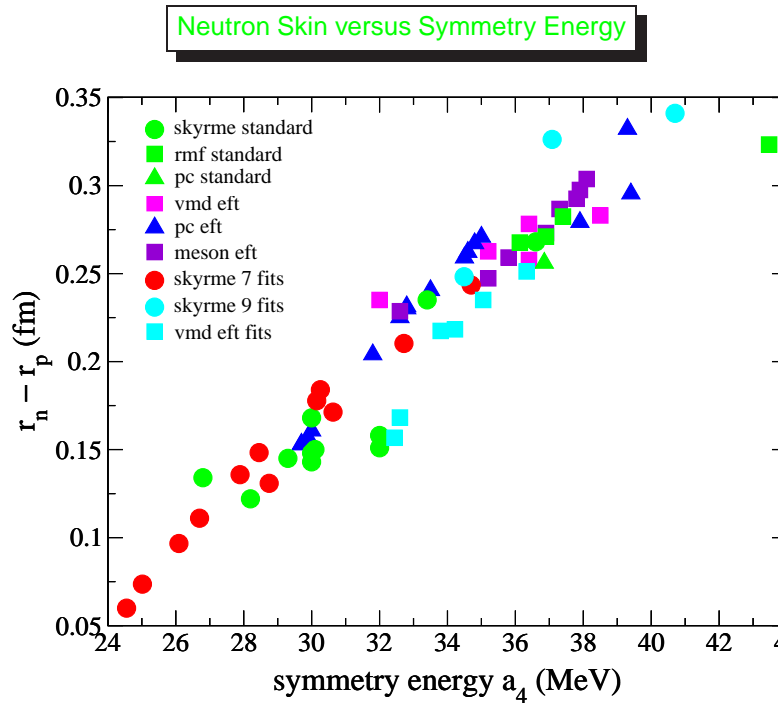
Comparison

- LOBF
- LOBHF+ 3BF
- Variational
- SCGF
- χ PT with pions
- DBHF
- RMF



Symmetry energy and neutron skin

Furnstahl et al pointed out correlation between a_4 , $p_0 = \rho_s^2 S'(\rho_s)$, and neutron skin, $\Delta R = R_n - R_p$ in ^{208}Pb



$$E(\rho, \alpha) = E(\rho, 0) + S_2(\rho)\alpha^2 + S_4(\rho)\alpha^4 + \dots \quad \alpha \equiv \frac{N-Z}{A}$$

$$E(\rho, 0) = -a_v + \frac{K_0}{18\rho_0^2}(\rho - \rho_0)^2 + \dots$$

$$S_2(\rho) = a_4 + \frac{p_0}{\rho_0^2}(\rho - \rho_0) + \frac{\Delta K_0}{18\rho_0^2}(\rho - \rho_0)^2 + \dots$$

Can one understand the S vs ΔR relation?

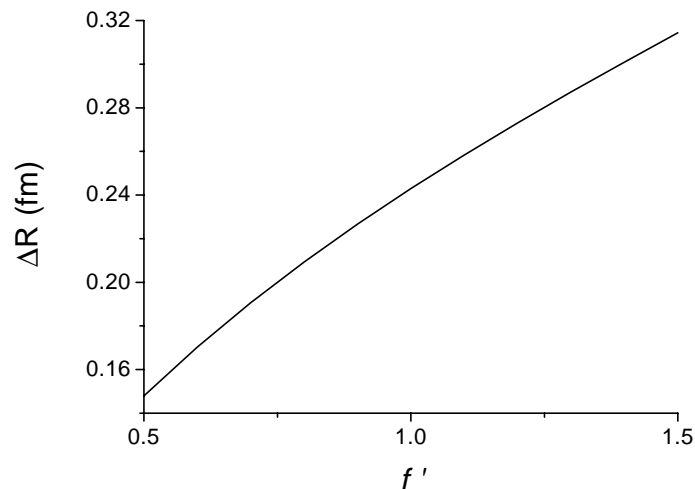
- Landau-Fermi liquid theory (thanks to V.Rodin)

$$V_{ph} = F + F' \tau_1 \cdot \tau_2 + G \sigma_1 \cdot \sigma_2 + G' \sigma_1 \cdot \sigma_2 \tau_1 \cdot \tau_2$$

In mean field approx $H_{MF} = T + V_0 + V_s \tau^3 + V_c$

- self-consistency: $V_s(r) = 2F'(\rho_n(r) - \rho_p(r))$

- Migdal relation: $a_4 = \frac{\epsilon_F}{3}(1 + 2F')$ relates SE with F'



- if F' density independent then $p_0 \sim a_4$

- SR for IVGMR: $r_n^2 - r_p^2 = F' \sum_s \frac{\langle 0 | N | s \rangle \langle s | \sum_i r_i^2 \tau_i^3 | 0 \rangle}{E_0 - E_s}$

Exp approaches to ΔR

- p,p' at 500-1000 MeV
- isovector Giant (spin) dipole resonance, IAS
- anti-protonic atoms
- parity violating electron scattering

Input of ΔR needed for

- weak charge Q_W
- deeply bound pionic atoms

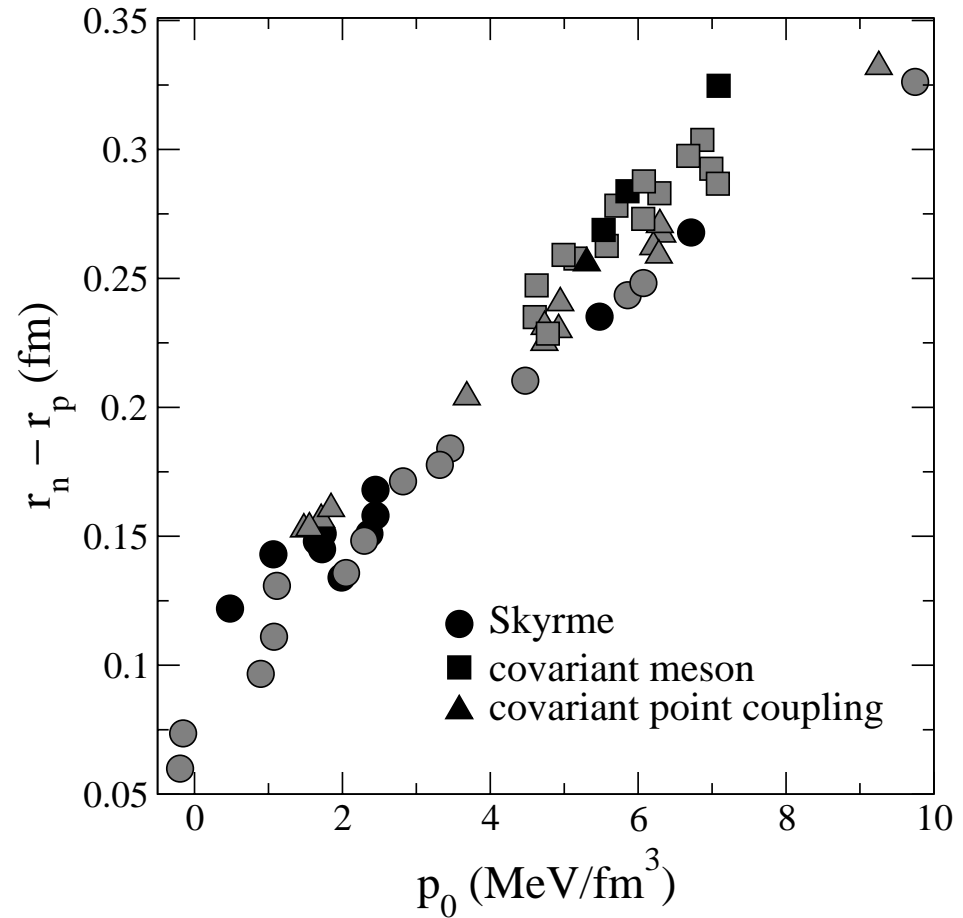
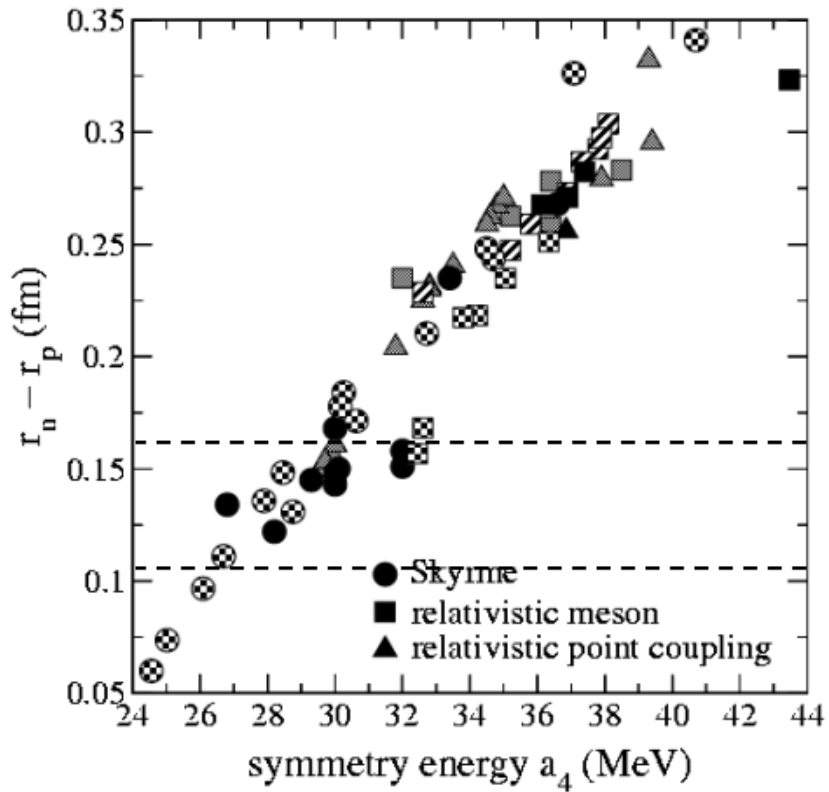
Present status ΔR in ^{208}Pb

| method | ΔR [fm] | Error [fm] | ref |
|--|-----------------|------------|----------|
| (\vec{p}, p') ($E = 0.5 - 1.0$ GeV) | 0.097 | 0.014 | Clark03 |
| nucleon scatt ($E = 40-200$ MeV) | 0.17 | | Karat02 |
| anti-protonic atoms | 0.15 | 0.02 | Trzcin01 |
| isovector giant dipole res. | 0.19 | 0.09 | Kras94 |
| giant dip. res. (unpubl.) | 0.14 | 0.05 | Kras02 |
| parity violating el sc | planned | 1% | Jlab |

- Exp consistent with $\Delta R \sim 0.10 - 0.15$ fm
- Clark et al: "Exp. seem to disagree with usual relativistic mean field models"
- New constraints might come from heavy ion reactions which probe higher densities

Symmetry energy and ΔR

In terms of mean field mc



Finite nuclei: LDA

- Effect of s.r.c. appears to be A -independent suggests use of LDA, see Van Neck et al, PRC51,1800 (1994)
- LDA Suppose one knows nm density $n^{nm}(k, k_F)$ as a function of $k_F = (\frac{3\pi^2}{2}\rho)^{1/3}$
- naive LDA $n_A(k) = \int dR n^{nm}(\rho(R), k)$ diverges for $k \rightarrow 0$
- improved LDA: apply to correlation part only
$$n_A(k) = n_A^{MF}(k) + \int dR \delta n_{<}(\rho(R), k) + \int dR \delta n_{>}(\rho(R), k),$$

where $n_A^{MF}(k) = \int dR \int dr \exp(ikr) \sum_h \phi_h^*(r_1) \phi_h(r_2)$

spectral function in LDA: $N=Z$

I refer to Van Neck et al, PRC 51 1800 (1994)

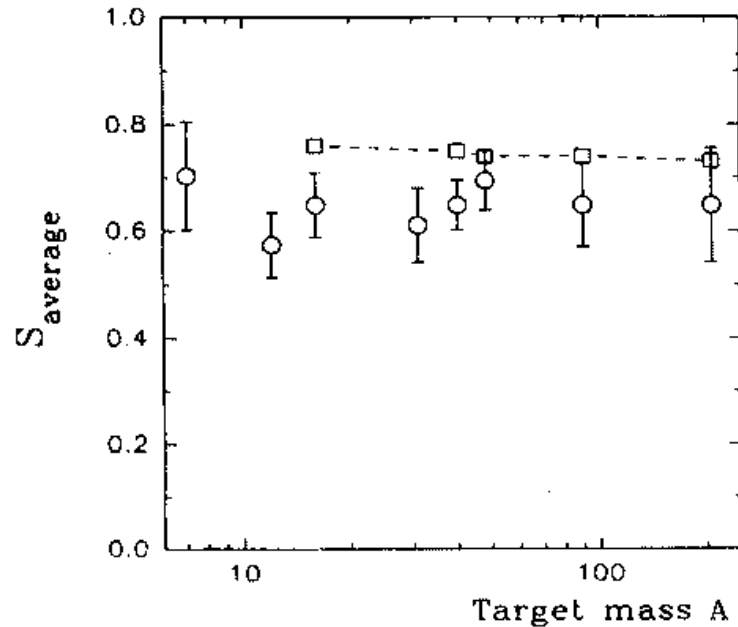
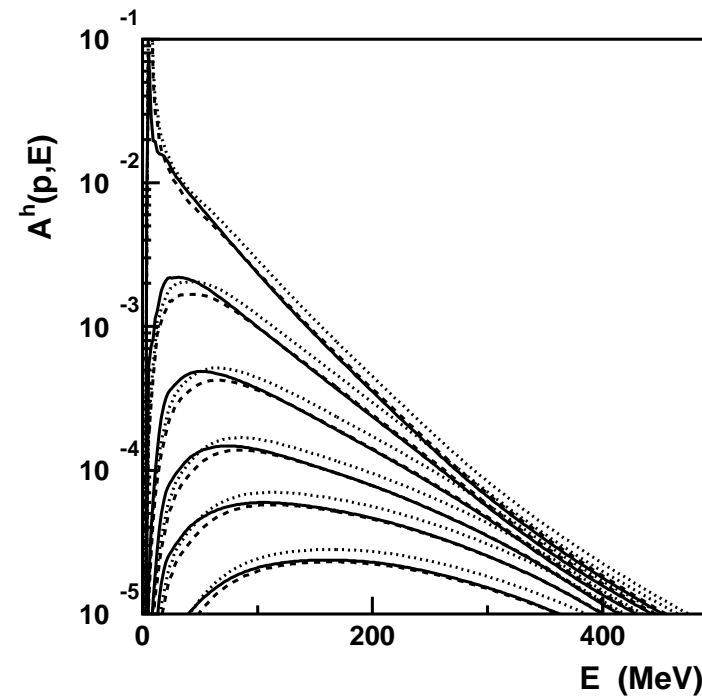
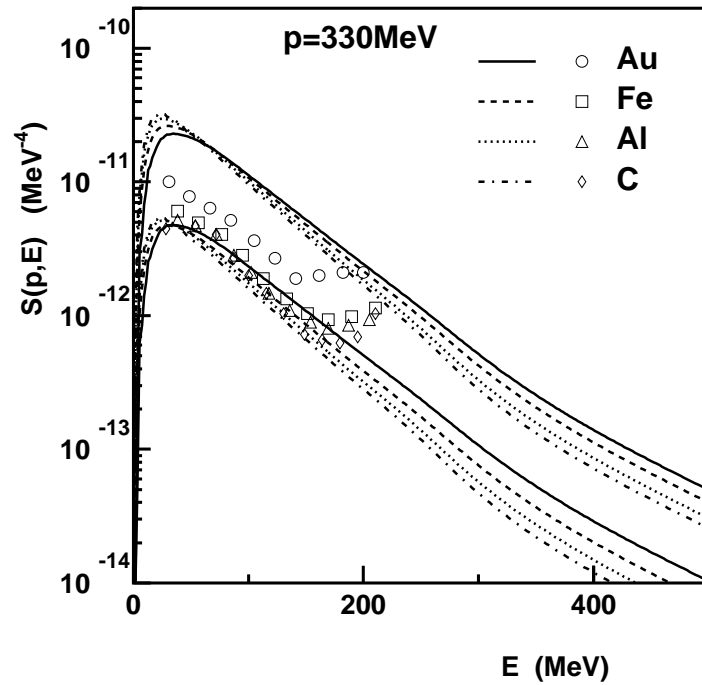


FIG. 5. Average spectroscopic factor for the valence hole

spectral functions in $N > Z$ finite nuclei



(from Bozek, nucl-th/0311046)