

# Nucleon knockout reactions on ${}^3,{}^4\text{He}$ induced by virtual photons

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- introduction
- exclusive  ${}^3,{}^4\text{He}(e, e'p)$
- semi-exclusive  ${}^4\text{He}(e, e'p)pnn$
- exclusive  ${}^3\text{He}(e, e'pp)$  and  ${}^3\text{He}(e, e'pn)$
- conclusions

# Introduction

To describe properties of few-body systems in terms of baryon-meson d.o.f. a microscopic model is needed.

$$\text{Hamiltonian: } H = -\frac{\hbar^2}{2m} \sum_i \nabla_i^2 + \sum_{i < j} v_{ij} \left( + \sum_{i < j < k} V_{ijk} \right)$$

Need a realistic NN-interaction:

$$v_{ij} = \sum_{p=1,N} v^p(r_{ij}) O_{ij}^p$$

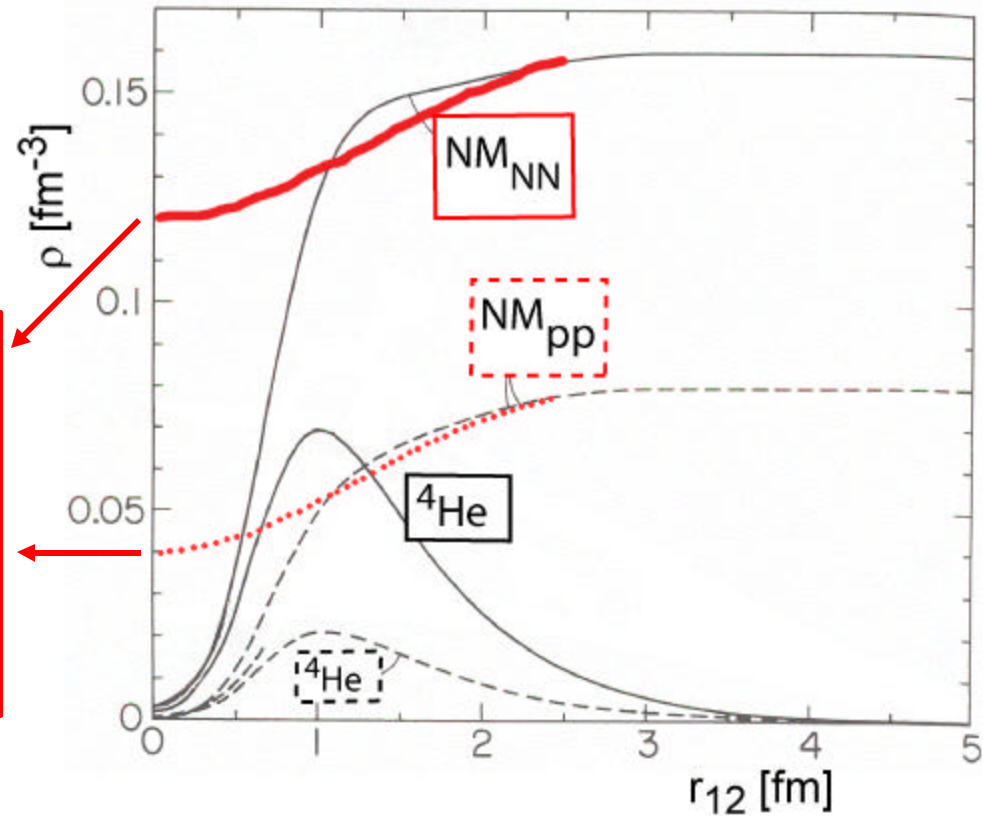
$$O_{ij}^p = 1, \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j, \boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j, (\boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j)(\boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j), S_{ij}, S_{ij}(\boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j), (\mathbf{L} \cdot \mathbf{S})_{ij}, \dots$$

Realistic nuclear forces induce spatial, tensor and spin-(iso)spin correlations, some of which are only known phenomenologically.

# two-nucleon density $\rho(\mathbf{r}_1, \mathbf{r}_2)$

—  $\rho_{NN}$   
- - -  $\rho_{pp}$

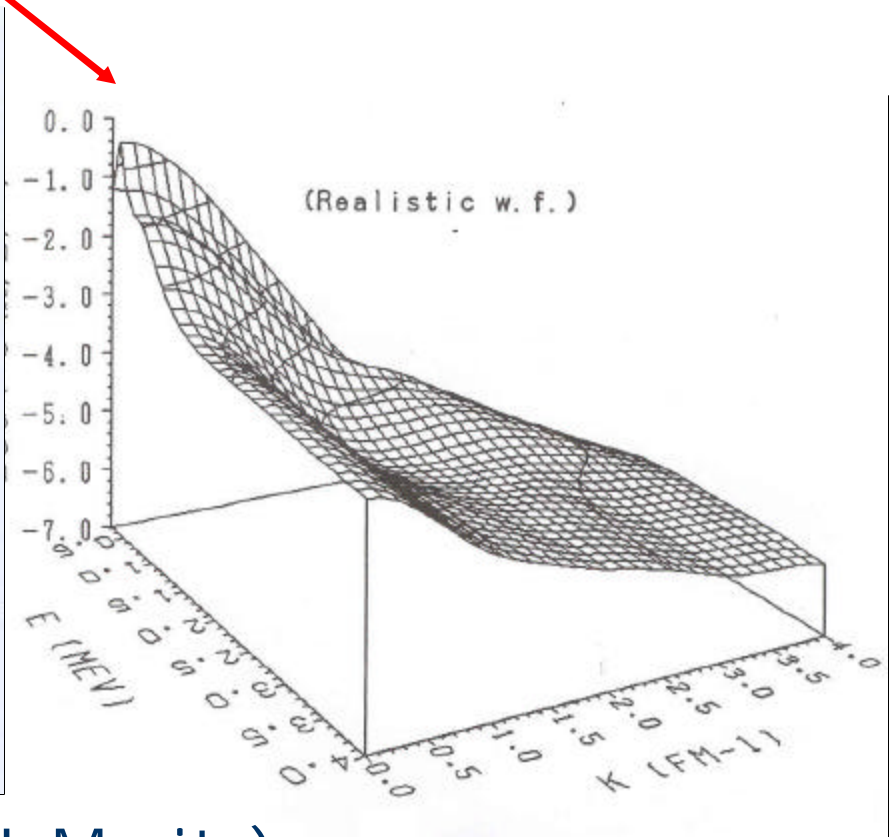
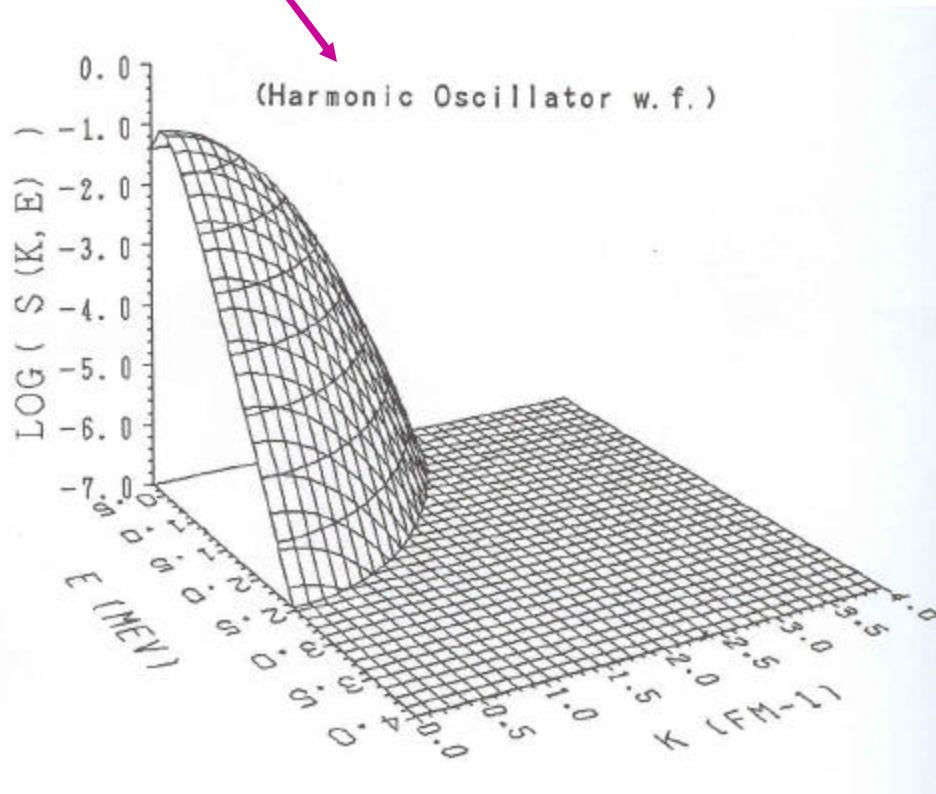
Slater  
determinant  
wave  
functions



Solve Schrödinger equation for 3N-system with Faddeev, VMC or GFMC.

Spectral function  $S(k,E)$

non-correlated and correlated proton spectral function  ${}^4\text{He}$



ATMS method (H. Morita)

Investigate high-momentum components with electro-induced single-nucleon knockout reaction

- clean probe, i.e.  $(\mathbf{q}, \omega)$  of the virtual photon
- well-defined final state
- light target nucleus:
  - \* initial and final state can be calculated precisely
  - \* disturbing processes can be calculated reliably

**exclusive  ${}^3, {}^4\text{He}(e, e'p)$ ,  ${}^2, {}^3\text{H}$  reaction**

Investigate strength at high missing momentum and high removal energies:

**semi-exclusive  ${}^4\text{He}(e, e'p)$**

Electron-induced two-nucleon knockout reactions to study initial and final-state correlations:

**${}^3\text{He}(e, e'pp)$  and  ${}^3\text{He}(e, e'pn)$**

# (e,e'p) reaction

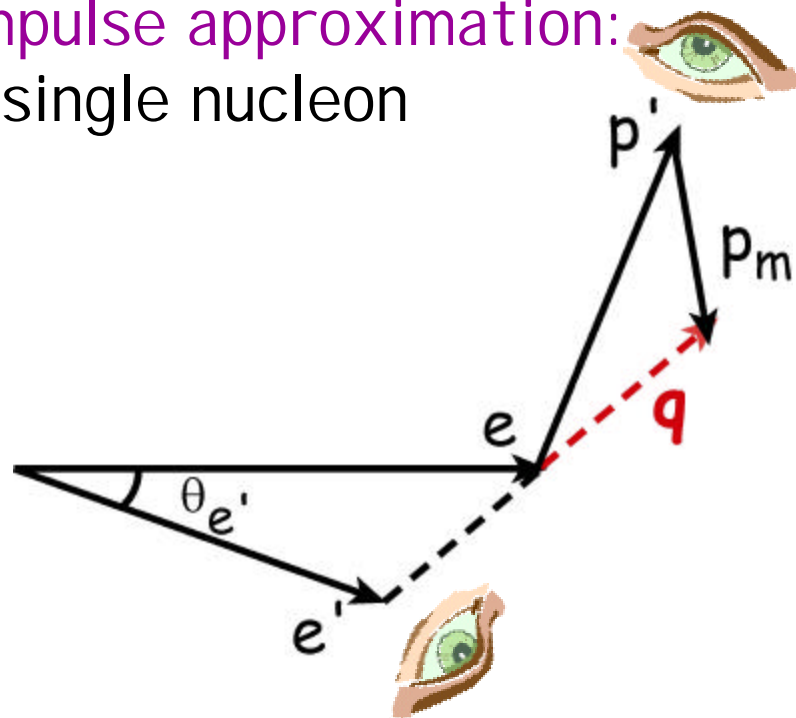
concept of (e,e'p) reaction in impulse approximation:  
virtual photon interacts with a single nucleon  
that subsequently is ejected

$$q = e - e'$$

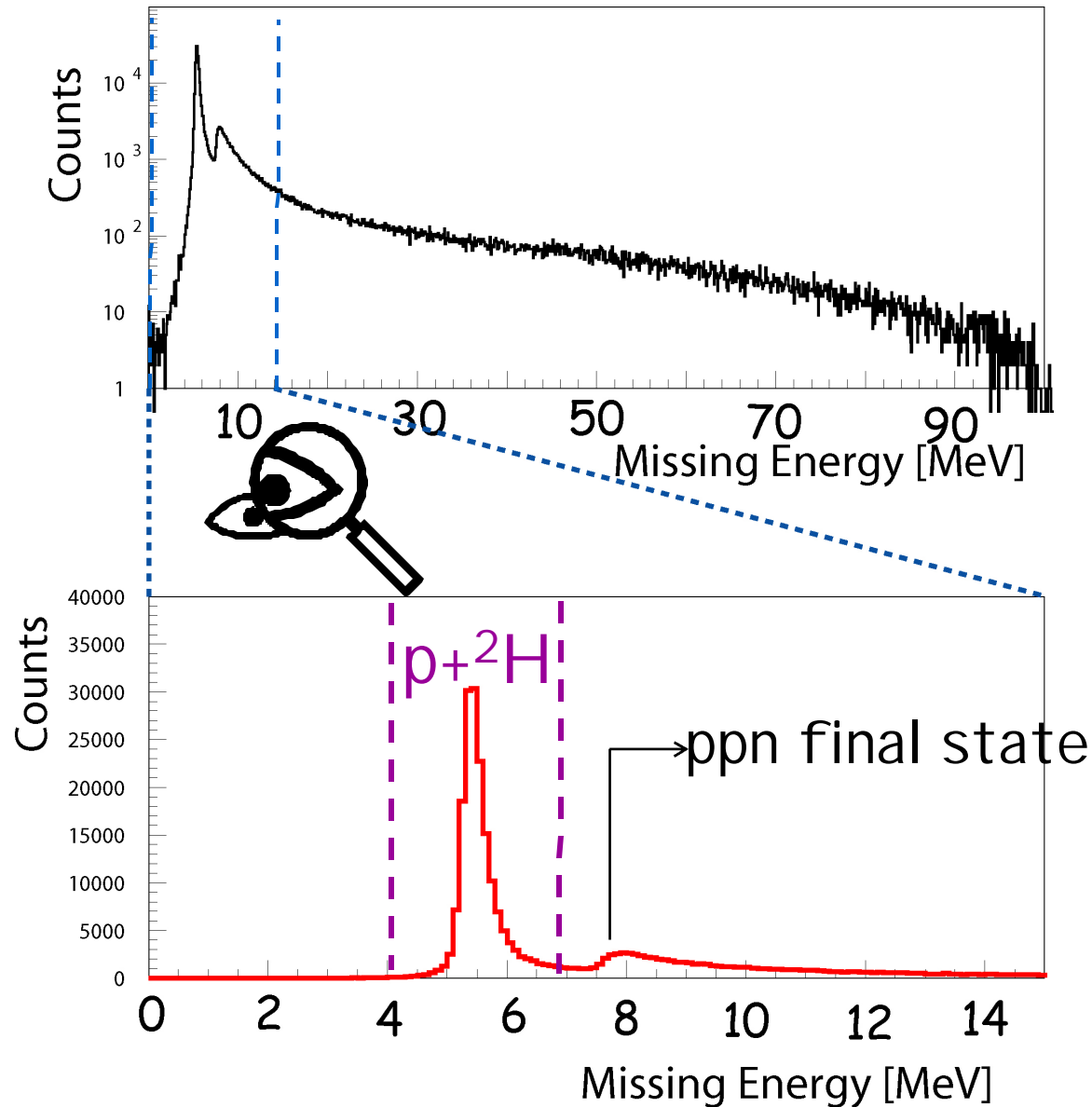
$$\vec{q} = \vec{e} - \vec{e}'$$

$$\vec{p}_m = \vec{q} - \vec{p}'$$

$$E_m = e - e' - T_p - T_{A-1}$$



# example $E_m$ -spectrum of ${}^3\text{He}(e,e'p)$ reaction



# ${}^3\text{He}(e, e'p){}^2\text{H}$ reaction

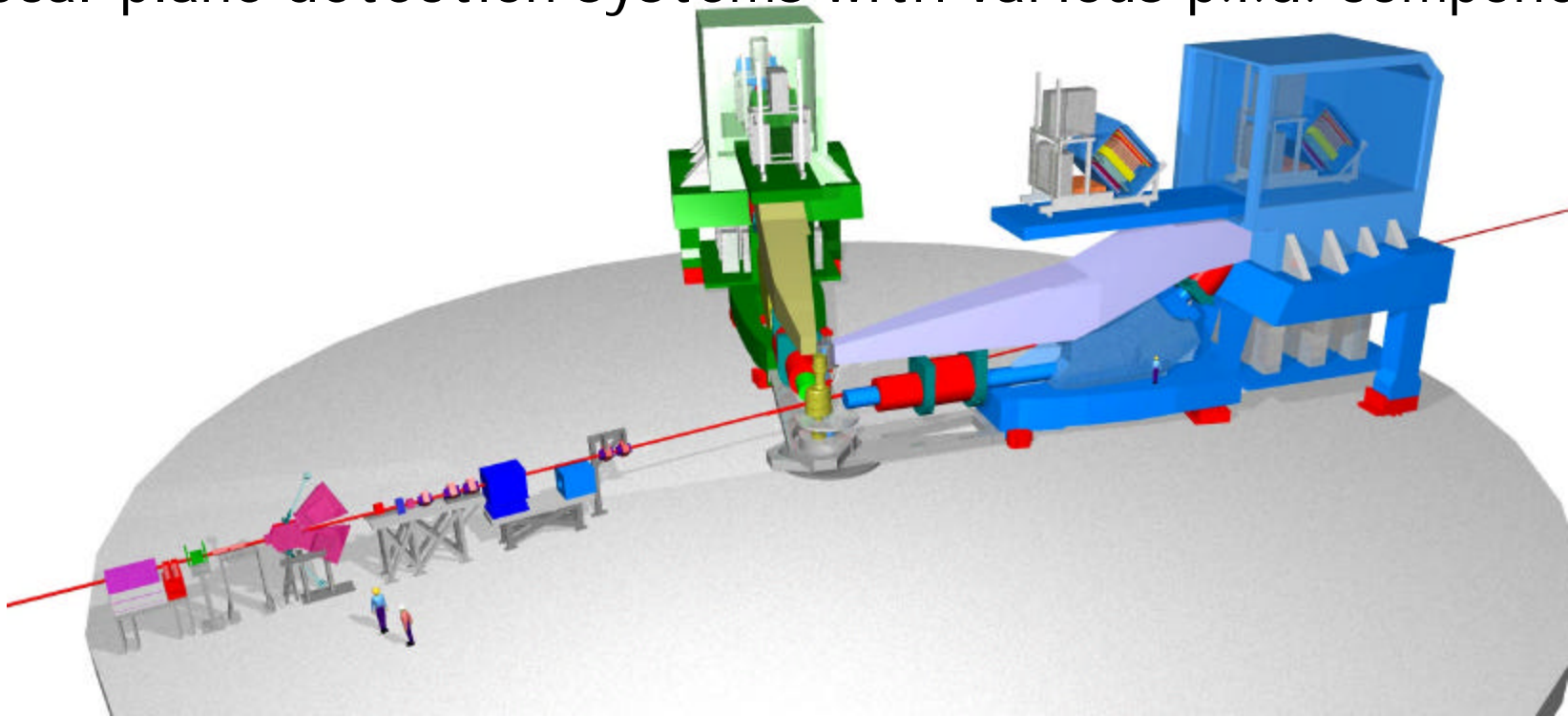
Hall-A of Jefferson Lab:

6 GeV 100% duty factor electron beam

high-power cryogenic targets

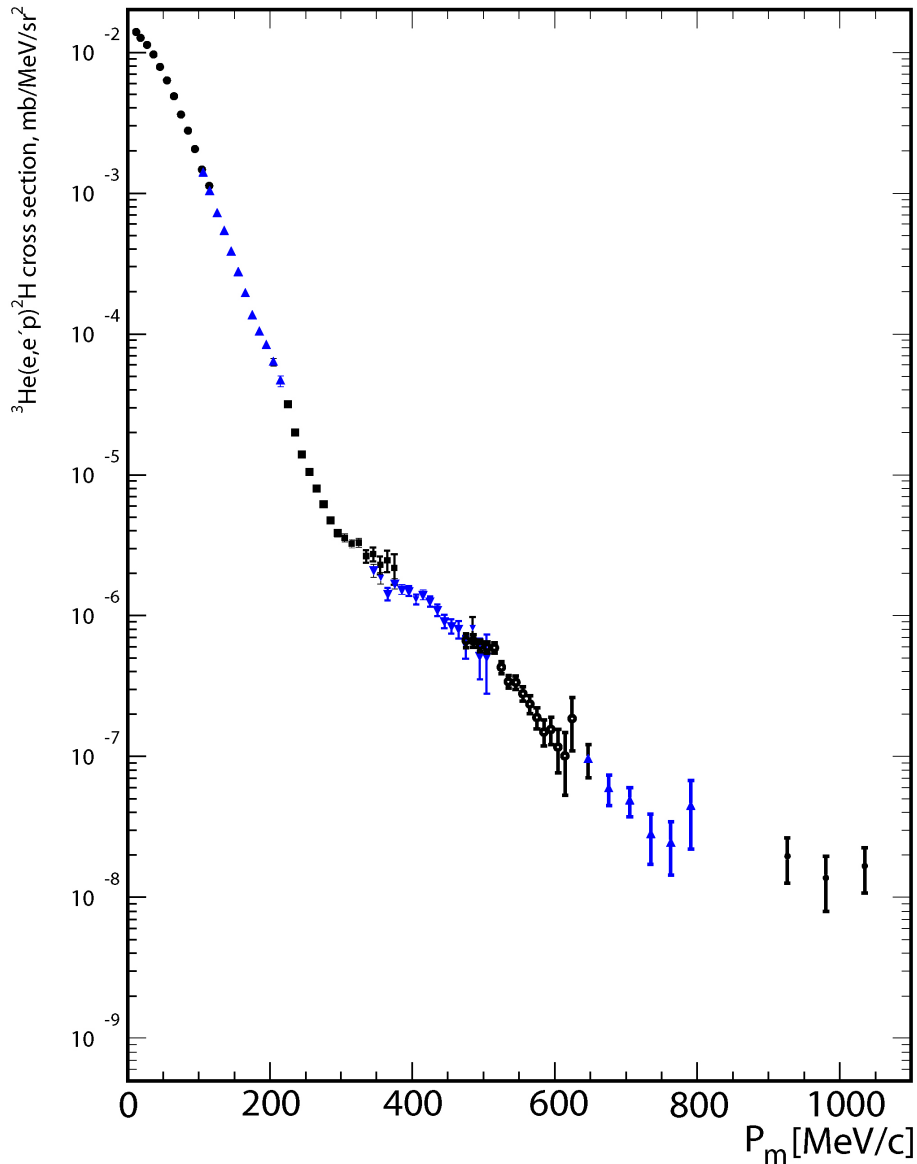
two high-resolution magnetic spectrometers

focal-plane detection systems with various p.i.d. components





# $p_m$ -distribution of ${}^3\text{He}(e,e'p){}^2\text{H}$



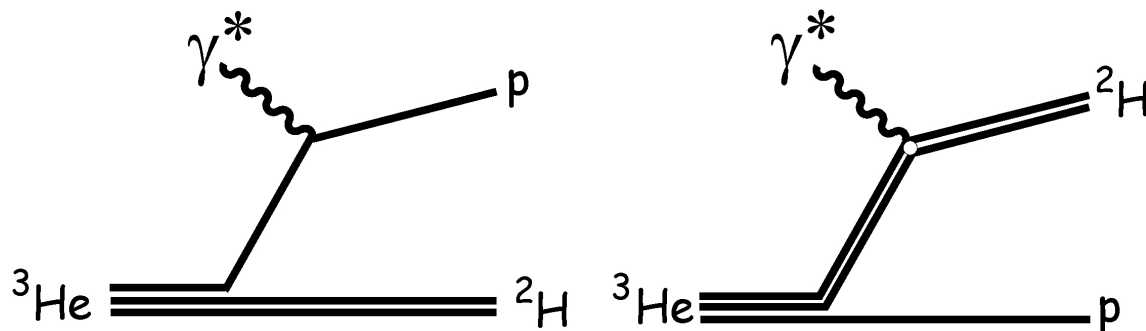
experiment E89-044  
(MIT-thesis of  
Marat Rvachev (2003))

$\omega=837$  MeV  
 $q=1500$  MeV/c  
 $Q^2=1.5$   $(\text{GeV/c})^2$

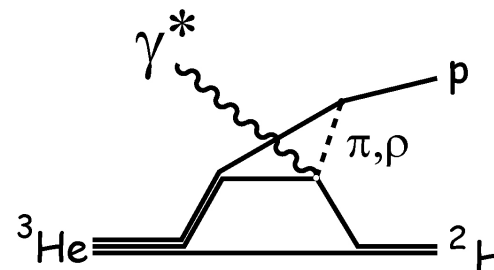
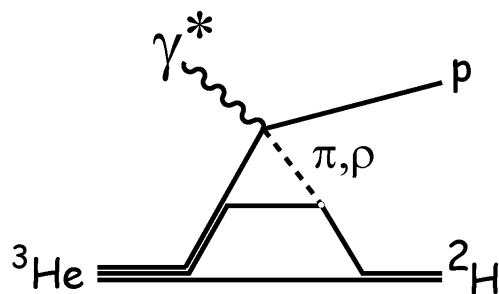
data up to  $p_m=1000$  MeV/c

# diagrammatic approach of Laget

PWIA

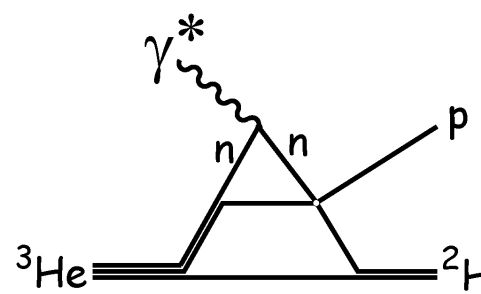
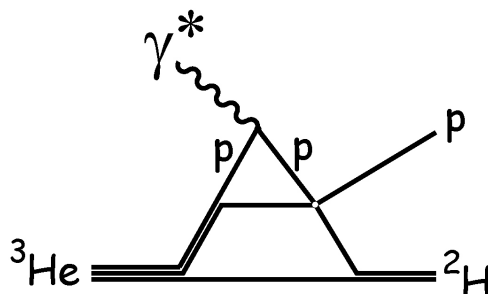


MEC



MEC

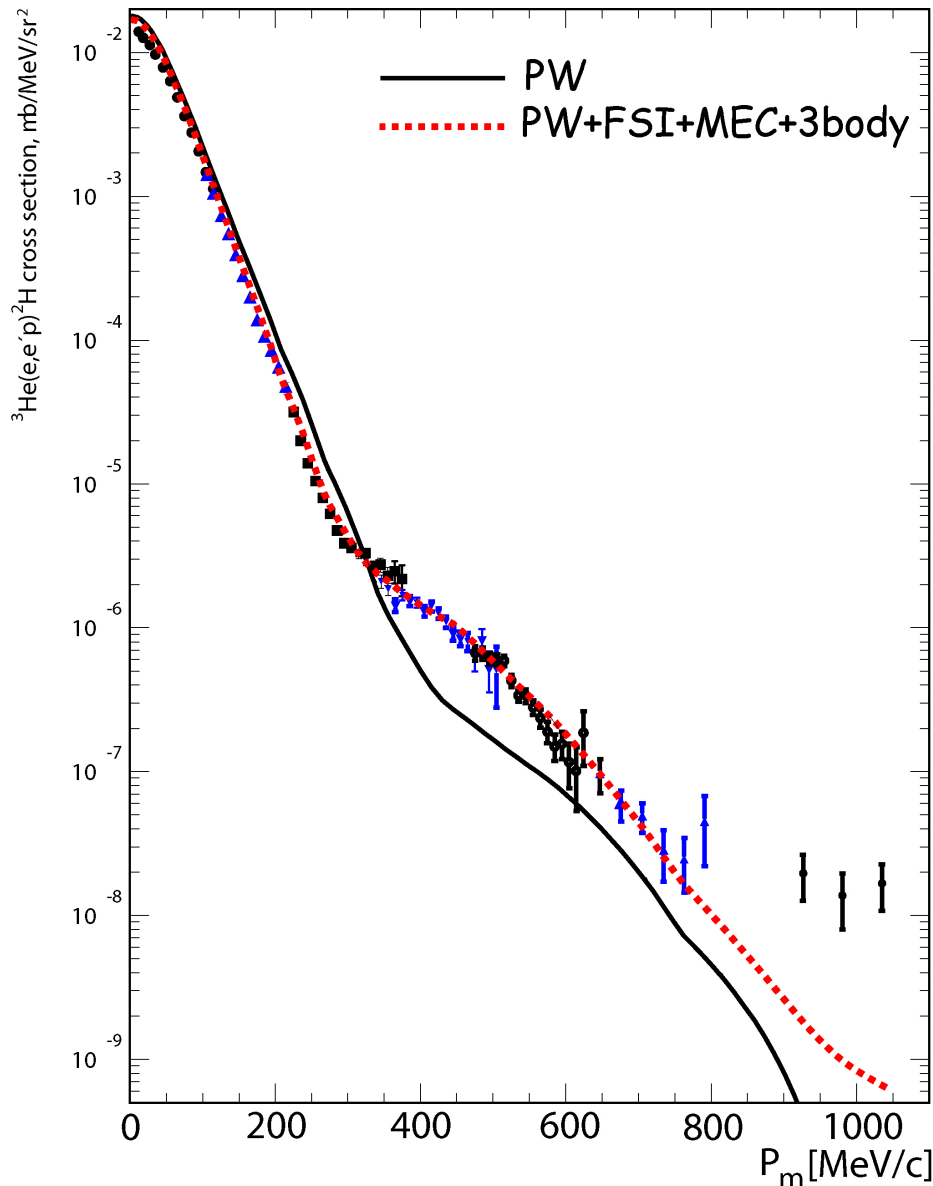
FSI



charge-exchange

+ three-body mechanisms

# $p_m$ -distribution of ${}^3\text{He}(e, e'p){}^2\text{H}$

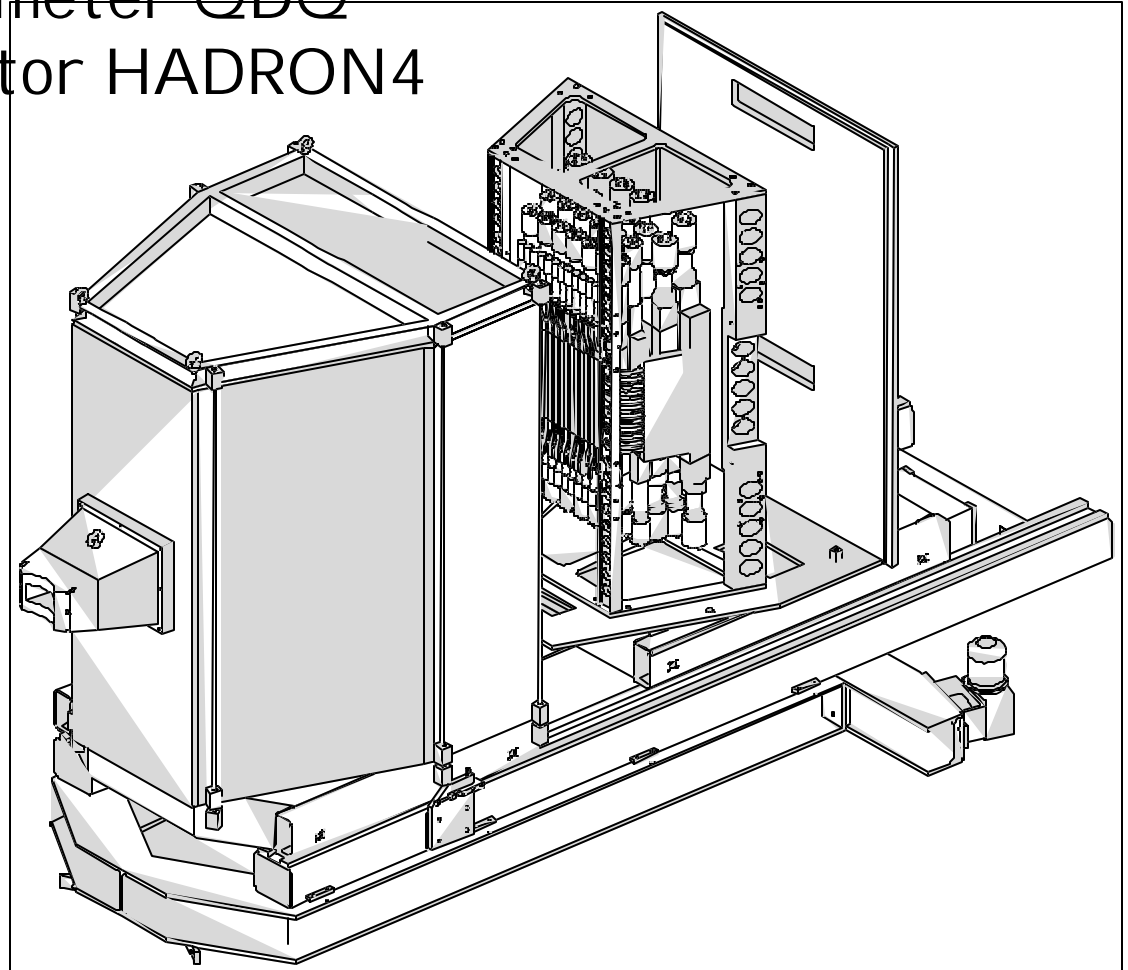


- large contributions from FSI and non-nucleonic currents at  $p_m > 400$  MeV/c
- excess strength at  $p_m > 800$
- relativistic effects ??
- breakdown of meson-baryon description ??

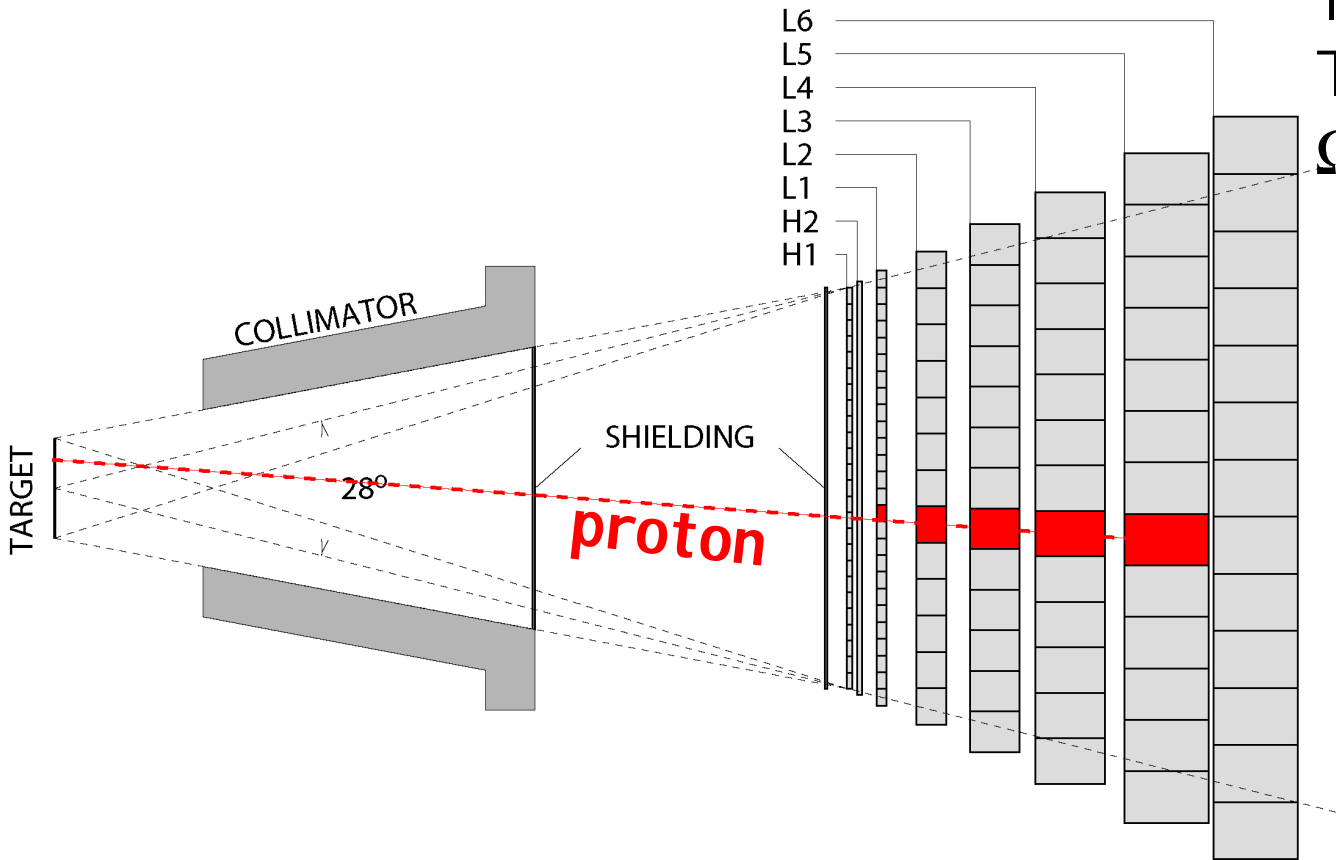
# $^4\text{He}(e, e'p)$ reaction @ NIKHEF

$e'$ : magnetic spectrometer QDO

$p'$ : segmented detector HADRON4



# HADRON4 for proton detection



94 scintillators  
134 PMs  
 $T_p = [67, 195]$  MeV  
 $\Omega = 550$  msr

determine:  
 $T_p$  and p.i.d.  
from generated  
light,  
angles with  
hodoscope

# ${}^4\text{He}(e, e'p)$ experiment @ NIKHEF

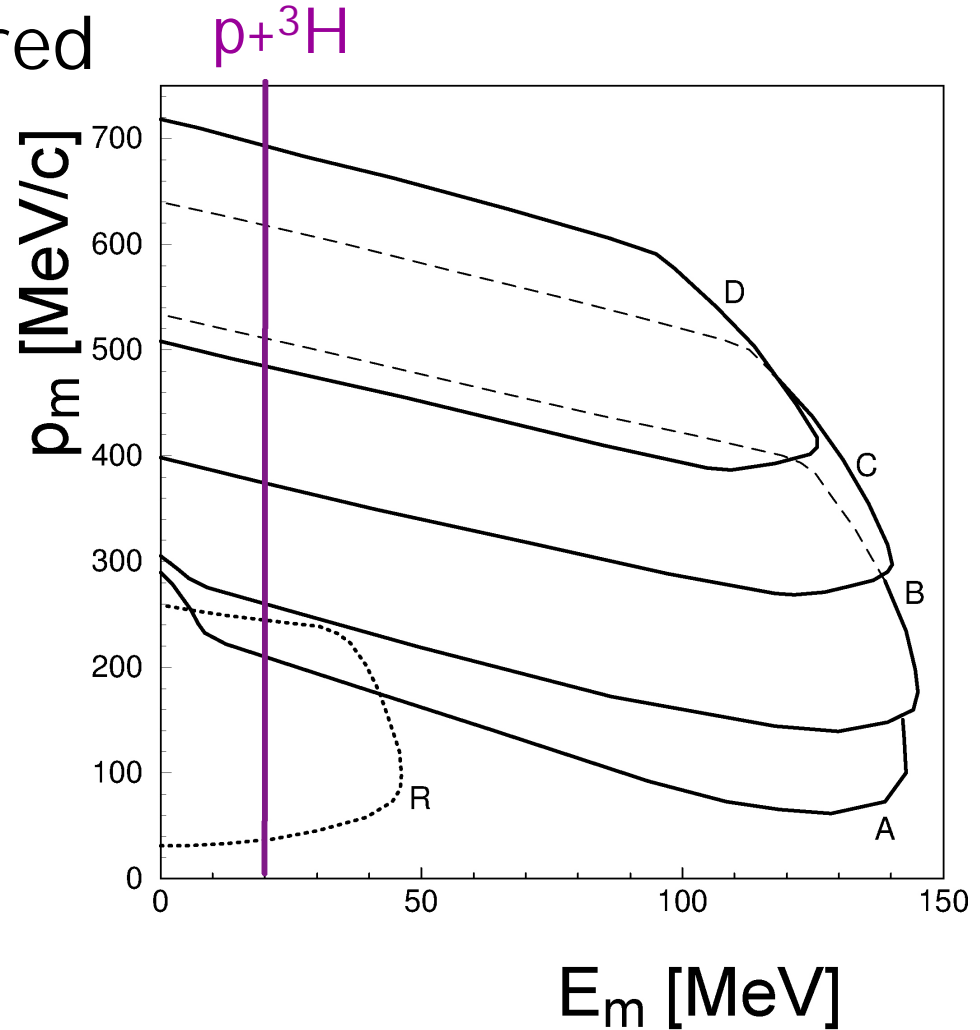
$(E_m, p_m)$  phase space covered

$$\omega = 215 \text{ MeV}$$

$$q = 400 \text{ MeV}/c$$

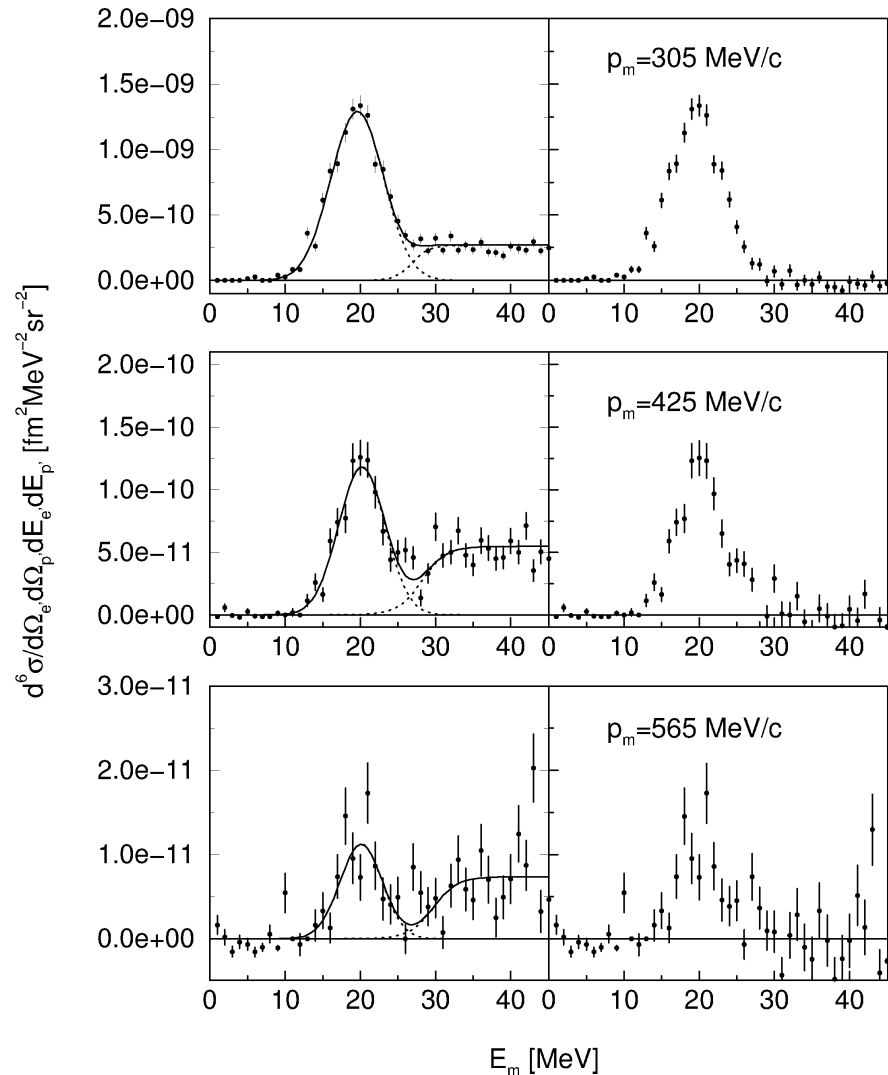
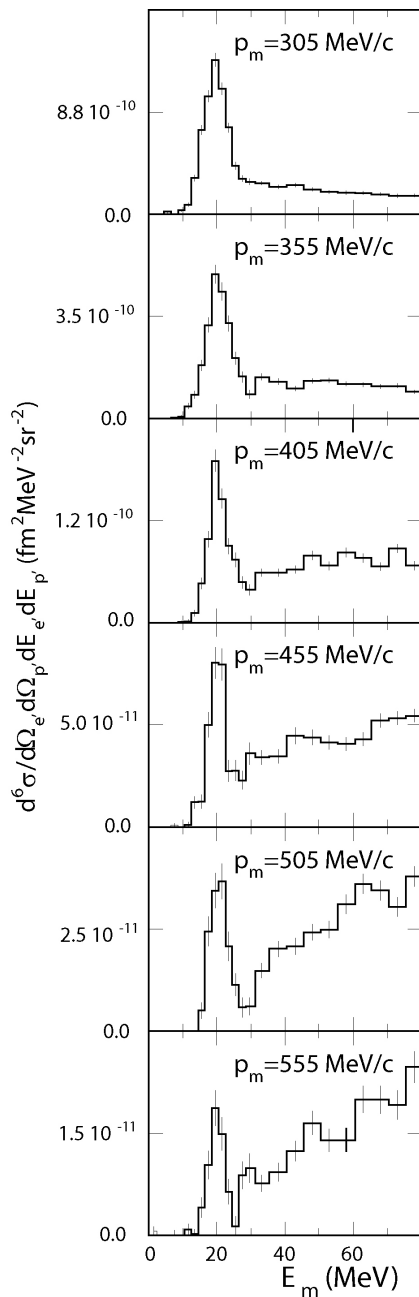
$$Q^2 = 0.11 \text{ (GeV}/c)^2$$

zero predicted in  
pT-momentum distribution  
around  $p_m$  400-500 MeV/c

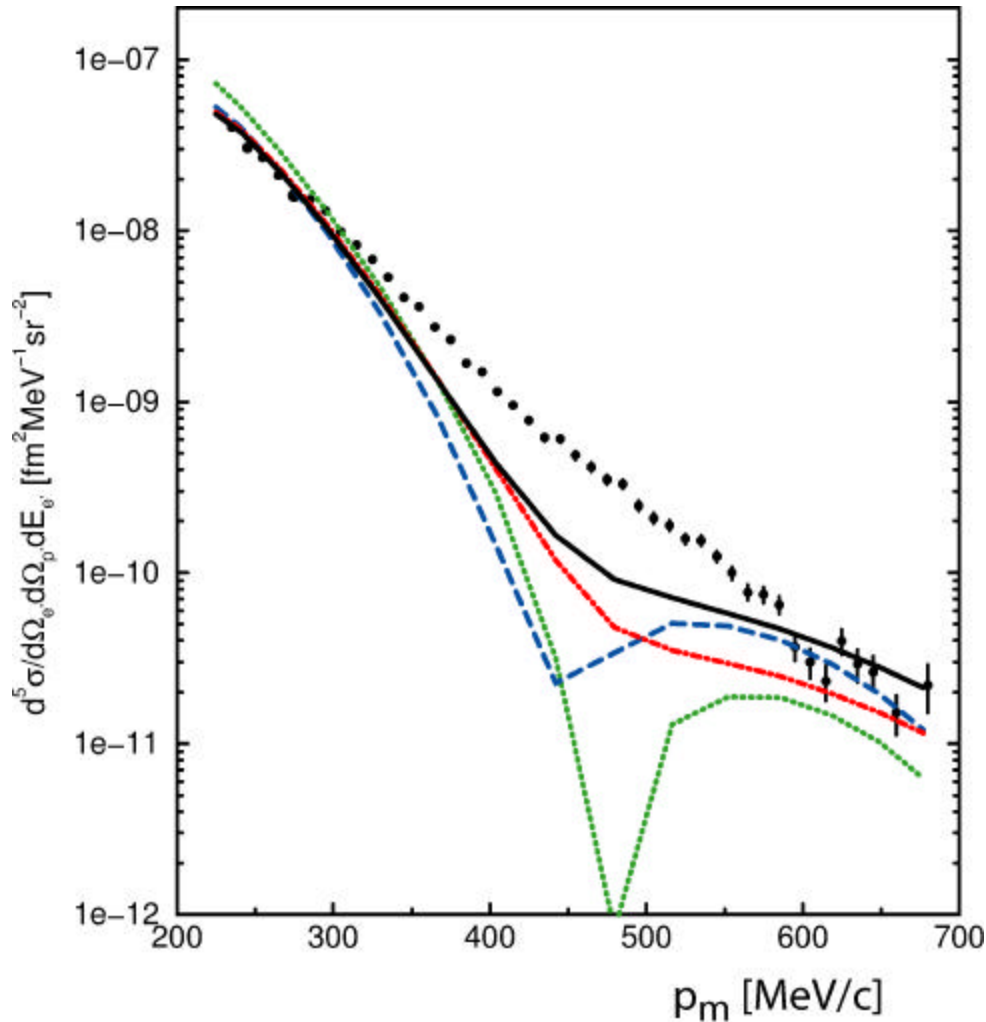


# $E_m$ -spectra of ${}^4\text{He}(e, e'p)$

selection of pT-breakup



# $p_m$ -dependence of ${}^4\text{He}(e, e'p){}^3\text{H}$



no zero/minimum  
observed in these data.

..... PWIA (v14+Urbana-VII)

----- +FSI

- · - · - · +FSI +MEC

———— full

More complete dataset of  
E97-111 of JLab

K. van Leeuwe et al. PRL **80** (1998) 2543

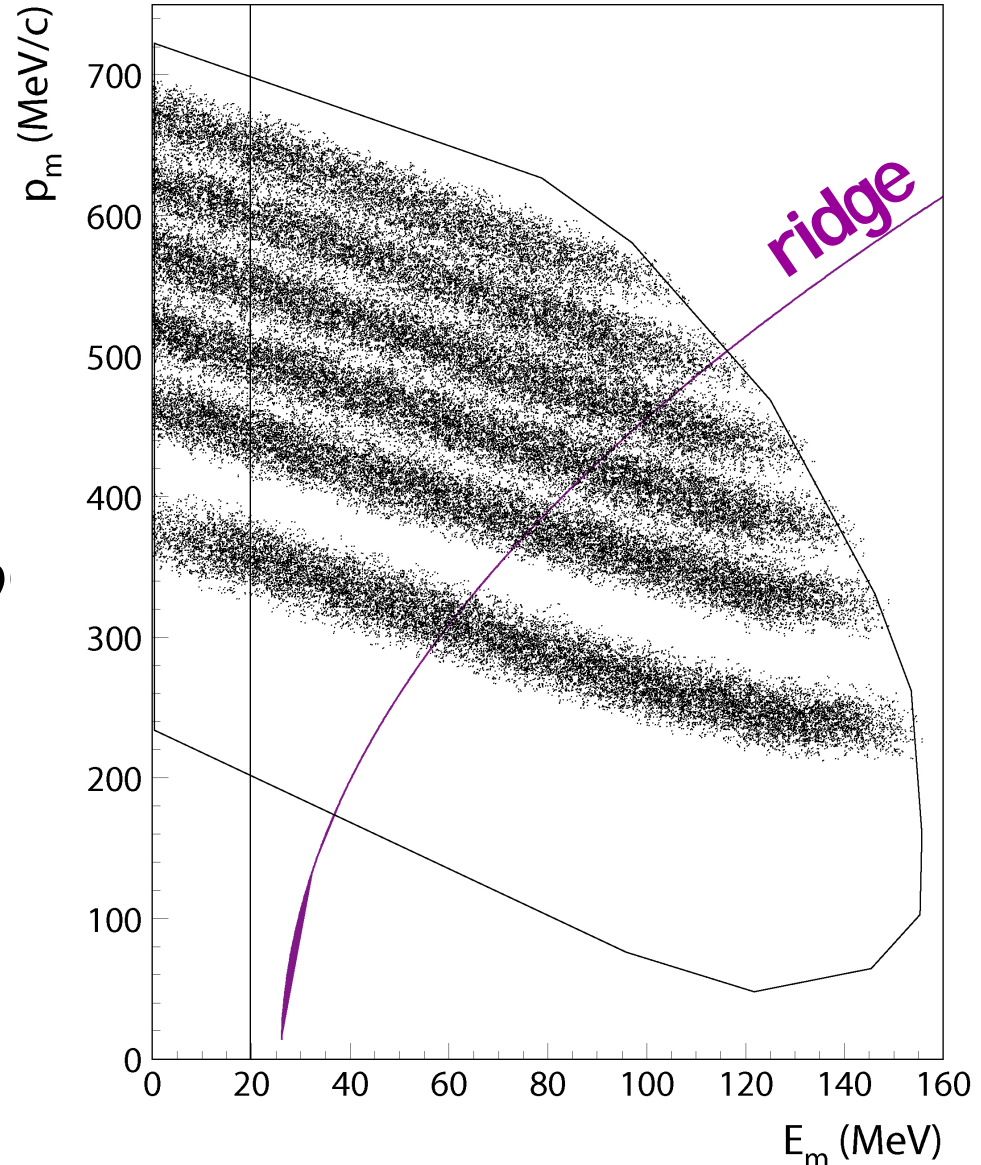


# semi-exclusive ${}^4\text{He}(e, e'p)$

scattering from a  
correlated pp/pn pair:

$$E_{m,\text{ridge}}(p_m) = E_{\text{thr}} + \frac{A-2}{A-1} \frac{p_m^2}{2M}$$

$5^\circ$  wide slices centered at  
 $\gamma_{pq} = 35^\circ, 50^\circ, 60^\circ, 70^\circ, 80^\circ, 89^\circ$

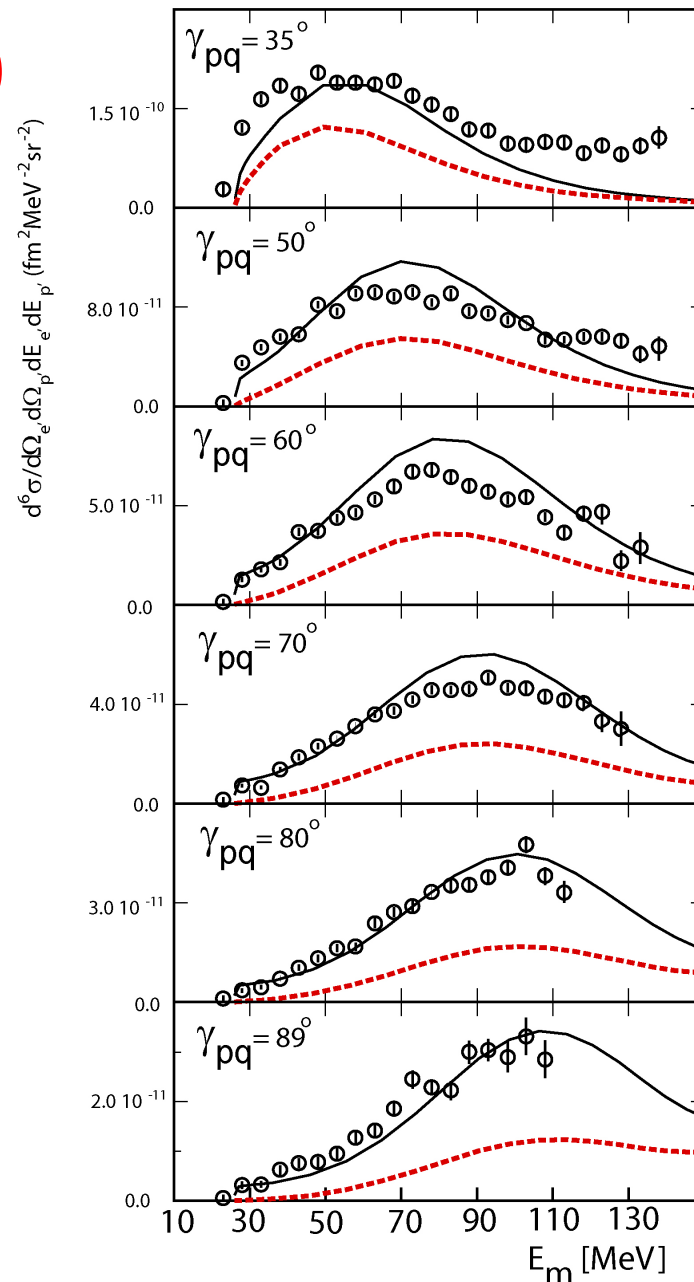


# $E_m$ -dependence( $\gamma_{pq}$ )

Laget

----- one-body +FSI

———— +MEC + IC



K. van Leeuwe et al.  
Phys. Lett. B523 (2001) 6

## summary of (e,e'p) on $^3,^4\text{He}$

- benchmark datasets for exclusive  $^3,^4\text{He}(e,e'p)$  up to very high  $p_m$ .
- signatures of scattering from a correlated nucleon-pair in semi-exclusive  $^4\text{He}(e,e'p)$ .
- advanced structure calculations of few-body systems.
- calculation of reaction dynamics needs to be improved.
- relativity needs to be included.

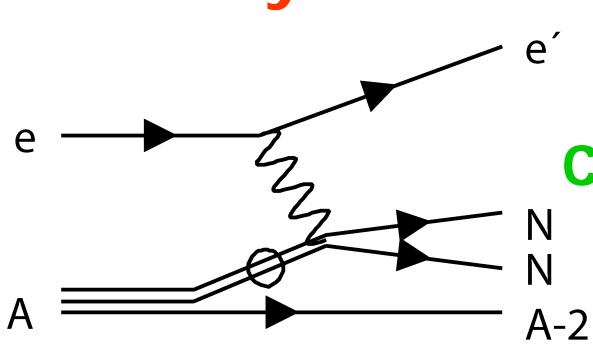
uncorrelated wave function  $\Rightarrow$  in PWIA  $\sigma(e, e'NN) = 0$

If one nucleon of a correlated pair gets hit by the virtual photon, both will presumably be emitted.

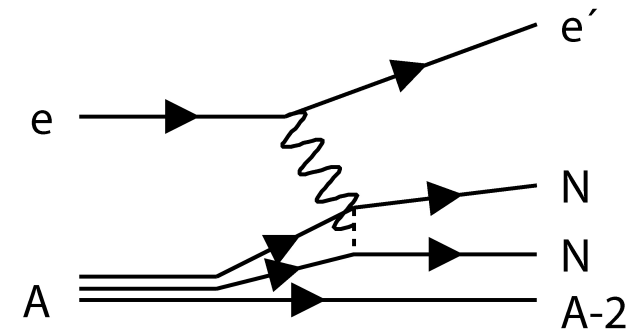
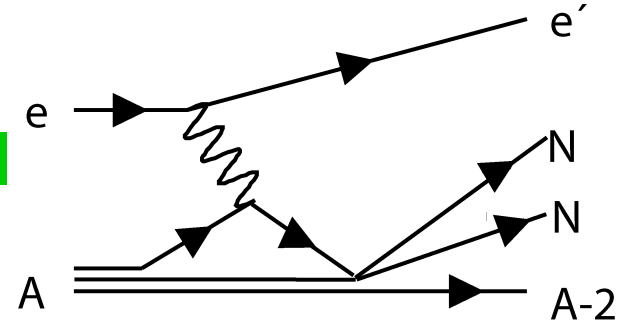
Two kind of processes in two-nucleon knockout:

- via **one-body current**  $j_1$ : pp and pn initial-state correlations
- via **two-body current**  $j_2$ : MECs and ICs

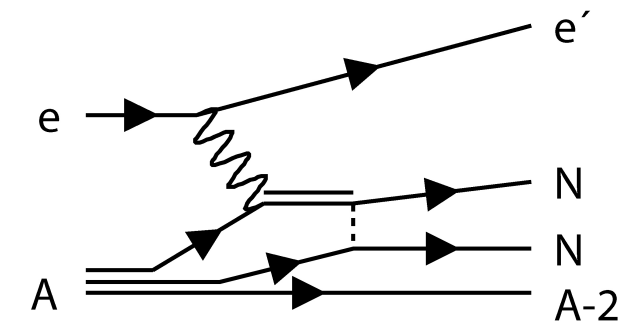
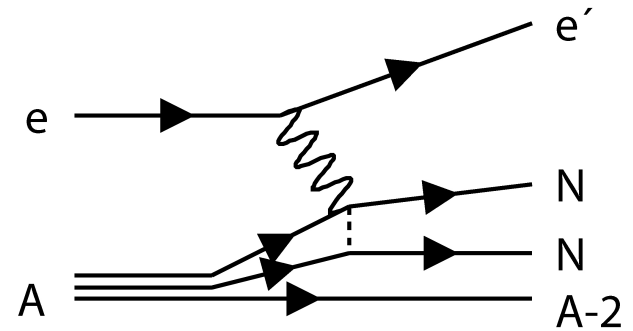
# Feynman diagrams of two-nucleon knockout



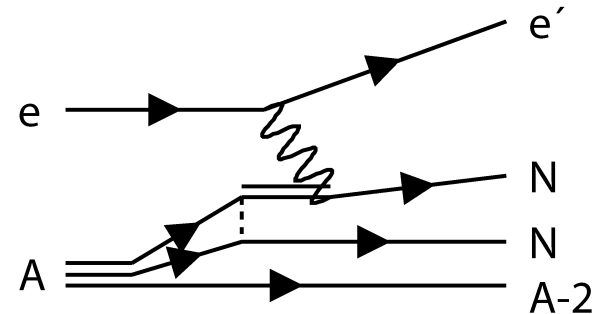
correlations  $j_1$  FSI



$j_2$   
MEC



$j_2$   
IC



Every process has its specific sensitivity to  $T$ ,  $q$ ,  $\omega$  and  $\epsilon$ .

Unpolarized (e,e'NN) cross section is a function of 6 nuclear structure functions.

**Disentangle the contributions  
via comparison of pp and pn-knockout  
under various kinematical conditions.**

$^3\text{He}$  is chosen because continuum Faddeev calculations are available for various realistic NN-potentials.

reconstruct  $p_m$  and  $E_m$  from measured  $e'$ ,  $p_1'$  and  $p_2'$

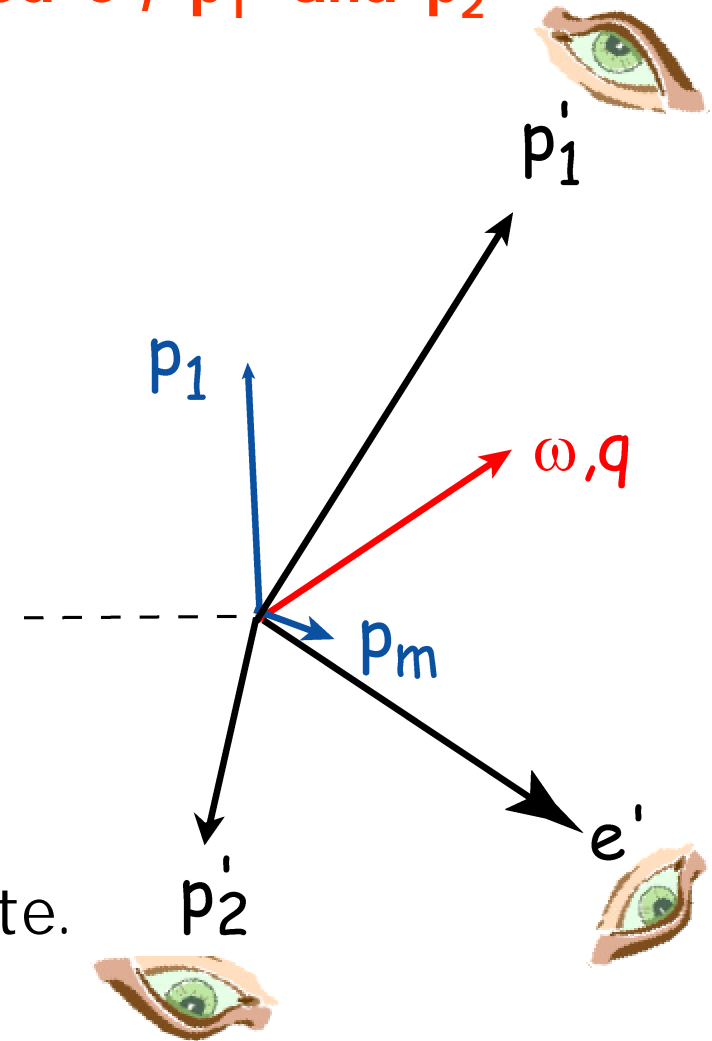
$$E_m = \omega - T_1 - T_2 - T_{(A-2)}$$

$E_m \sim E_{\text{exc}}$  of (A-2) system

$$\vec{p}_m = \vec{q} - \vec{p}_1' - \vec{p}_2'$$

$^3\text{He}$ : (A-2)=nucleon, so  $p_m$  is momentum of the (unobserved) nucleon in the final state.

in PWIA:  $-\mathbf{p}_m = \text{CoM momentum of the pair in the initial state.}$



**$^3\text{He}(e, e'pp)$  experiment performed at AmPS  
(NIKHEF, Amsterdam)**

D. Groep et al., PRL **83** (1999) 5443

D. Groep et al., PR C**63** (2000) 014005

**$^3\text{He}(e, e'pn)$  experiment performed at A1  
(MAMI, Mainz)**

*NIKHEF, Amsterdam, The Netherlands*

*Institut für Kernphysik, Mainz, Germany*

*Physikalisches Institut, Tübingen, Germany*

*University of Glasgow, Glasgow, Scotland*

**measurements at the same central values of  $(w, q)$**



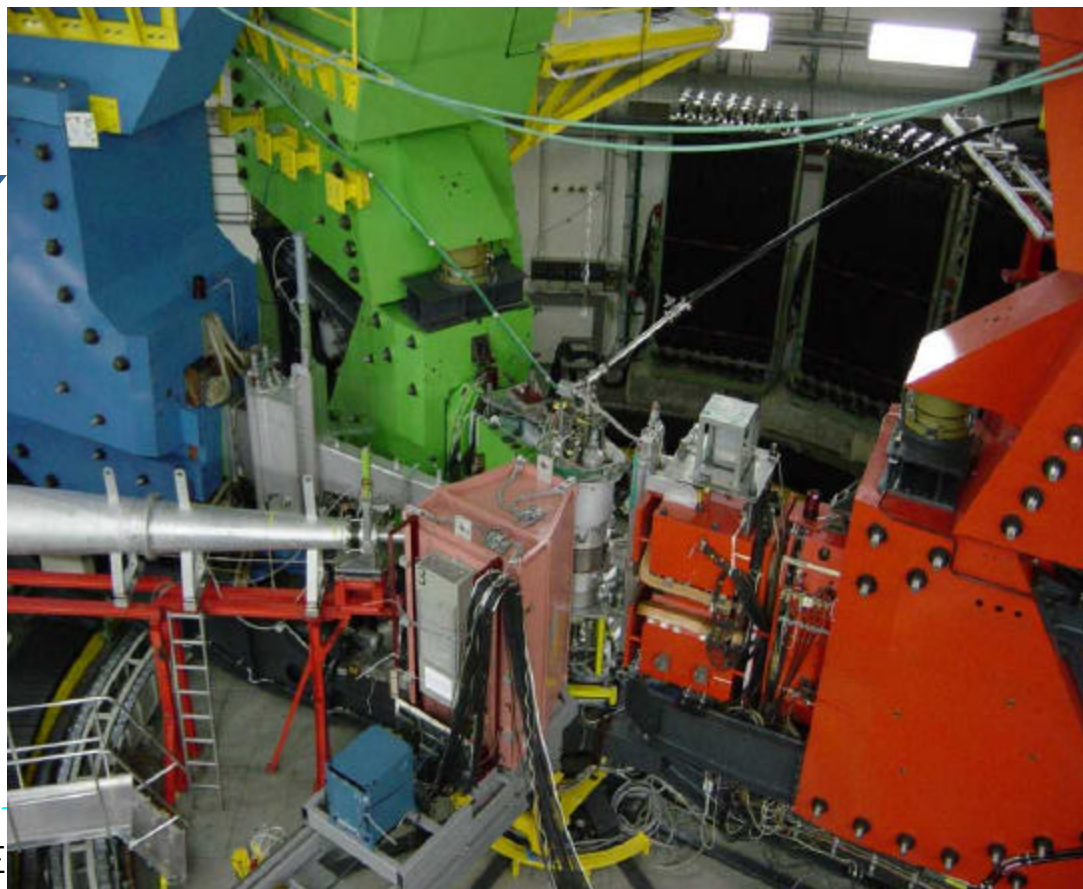
## experimental tools

- high duty-factor electron beam
- cryogenic target
- magnetic spectrometers for electron detection
- large solid angle scintillator detectors, with a high degree of segmentation, for proton & neutron detection

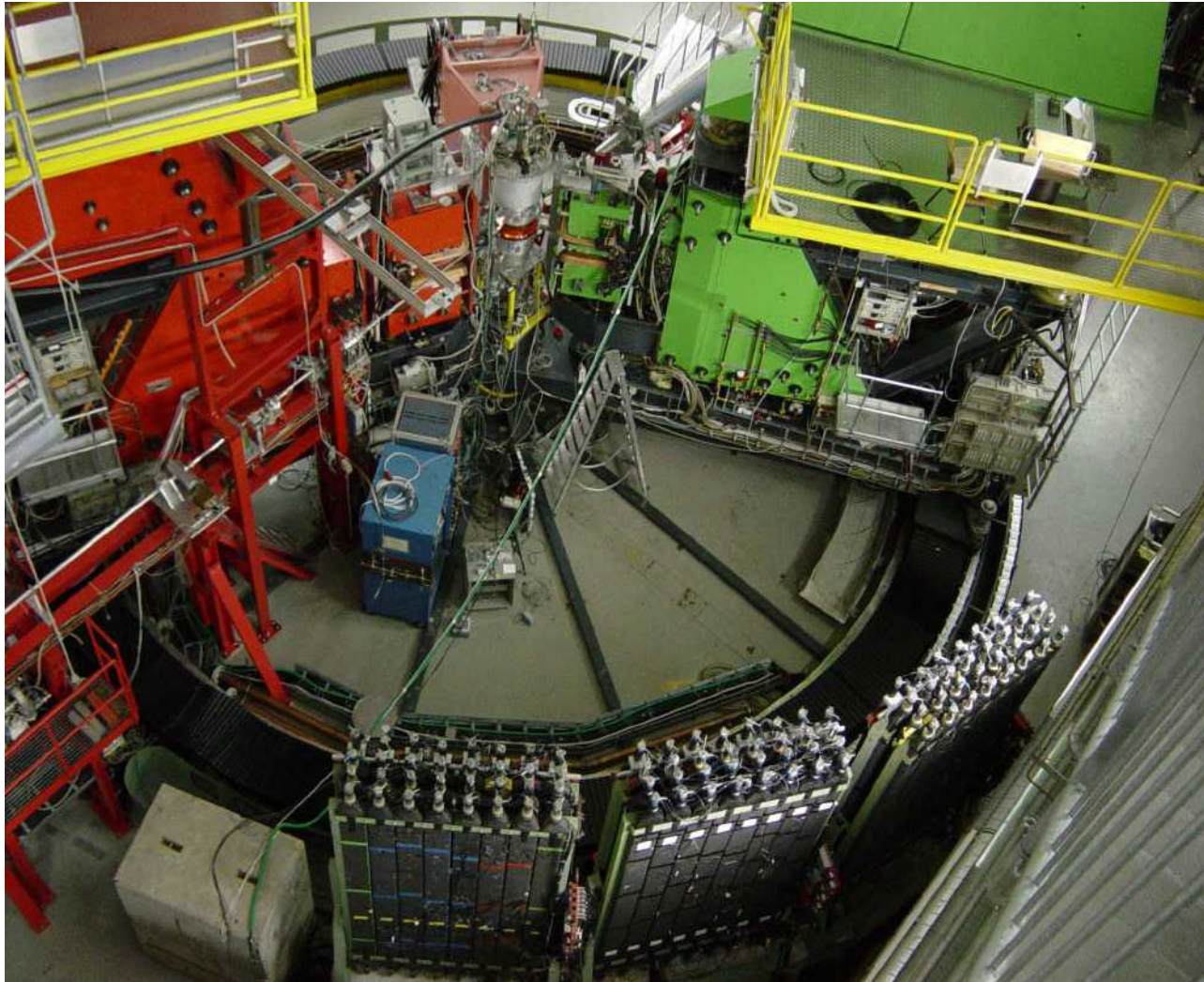
A1-hall of MAMI

spekB:  $\Omega=5.6$  msr  
 $\Delta p/p=15\%$

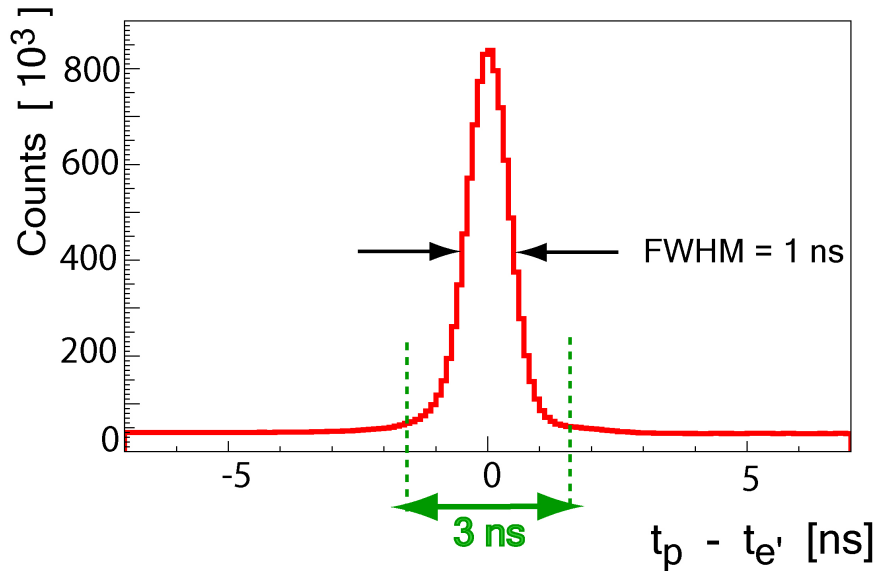
**Luminosity:**  
 $10^{36} \text{ cm}^{-2} \cdot \text{s}^{-1}$



# A1-hall at MAMI

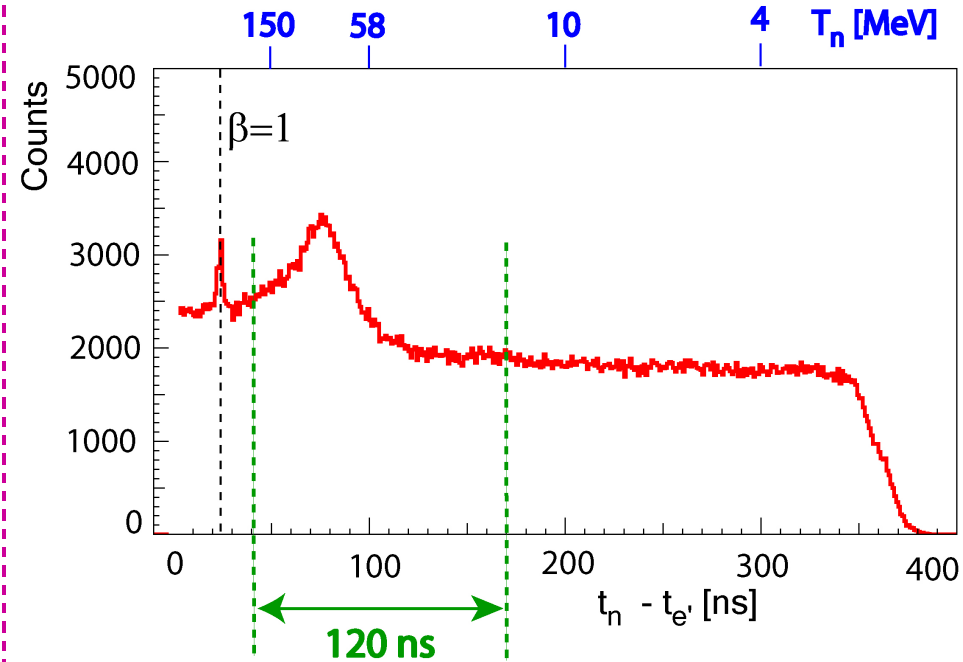


# electron-proton time difference spectrum



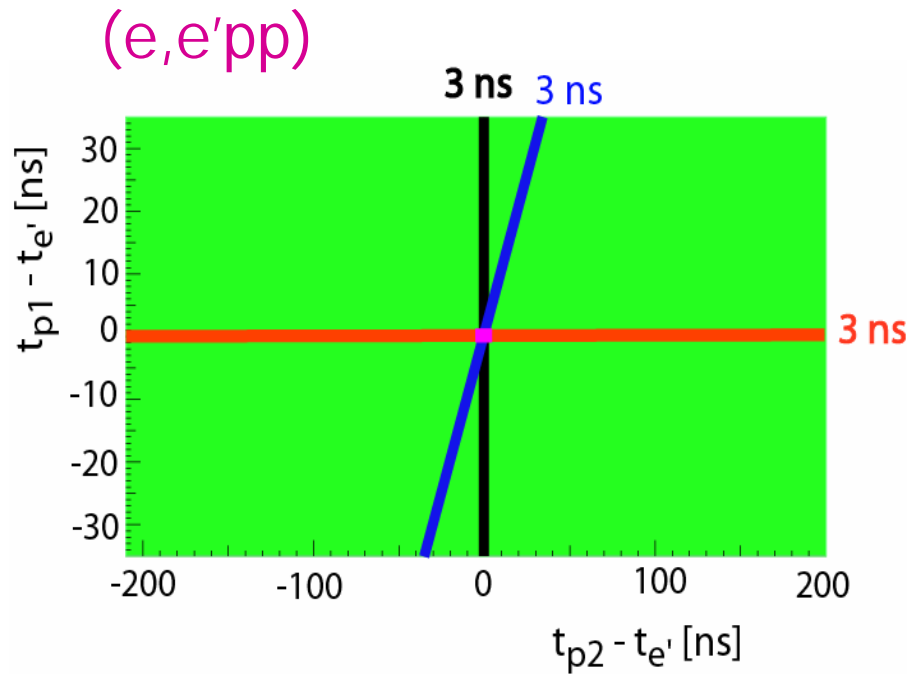
time and energy of proton are uncorrelated

# electron-neutron time difference spectrum



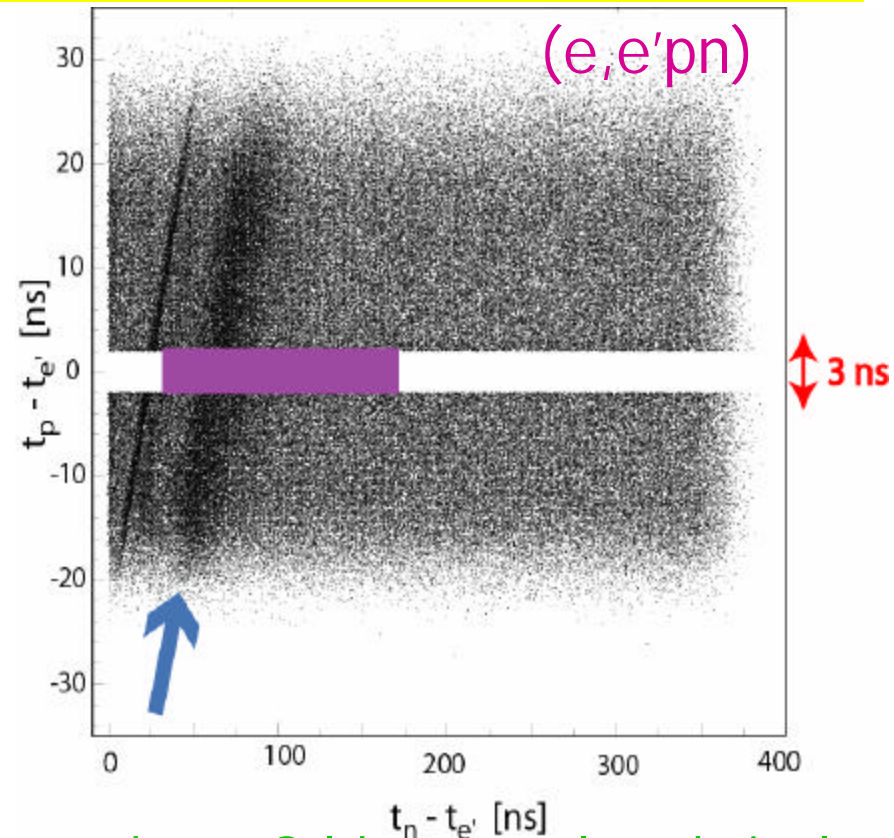
time and energy of neutron are correlated

# coincidence times and accidental subtraction procedures



- three-fold uncorrelated singles
- real-(e'p<sub>1</sub>) + accidental n (3 ns)
- real-(e'p<sub>2</sub>) + accidental p<sub>1</sub> (3 ns)
- real-(p<sub>1</sub>p<sub>2</sub>) + accidental e' (3 ns)

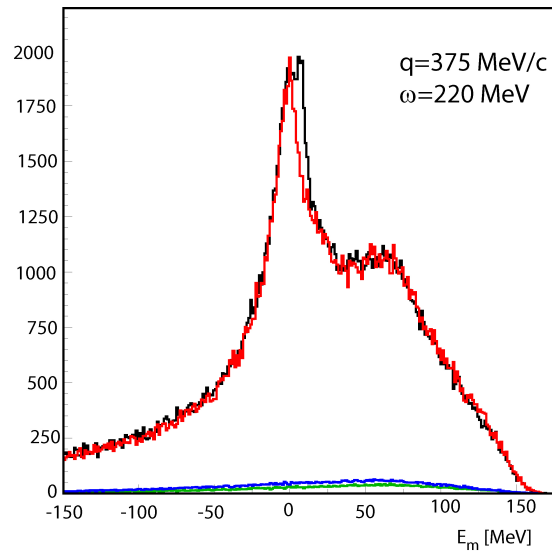
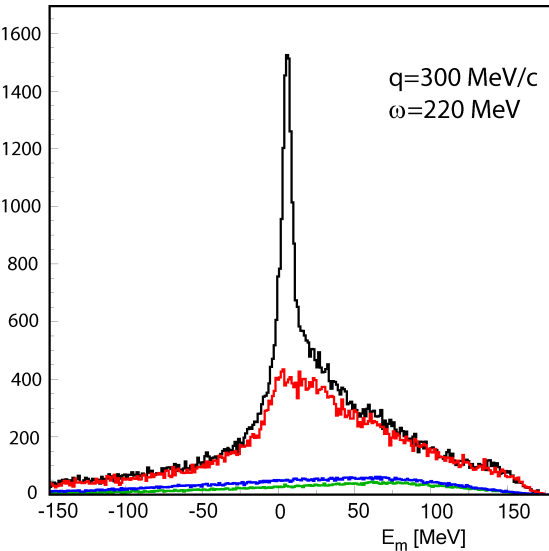
real (e,e'pp) in 3x3 ns<sup>2</sup>



- three-fold uncorrelated singles
- real-(e'p) + accidental n (3 ns)
- real-(e'n) + accidental p (120 ns)
- real-(pn) + accidental e' (120 ns)

real (e,e'pn) in 3x120 ns<sup>2</sup>

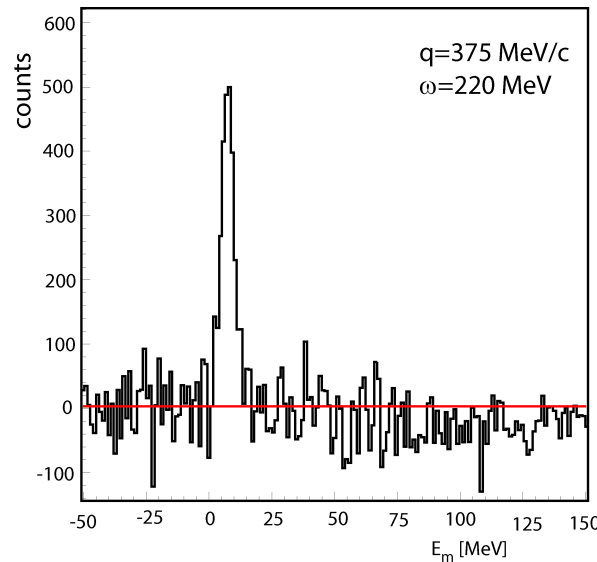
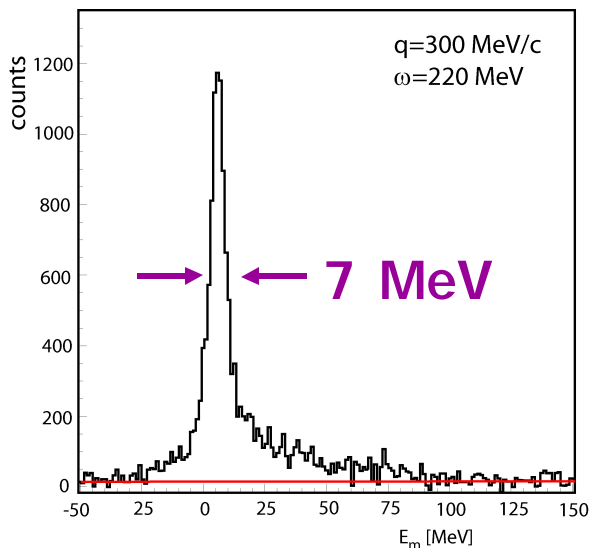
# $E_m$ spectra of ${}^3\text{He}(e, e'pn)$ @ MAMI



in  $3 \times 120 \text{ ns}^2$

- total (e'pn)
- 3x uncorrelated singles
- real-(e'n) + acc. p
- real-(pn) + acc. e'
- real-(e'p) + acc. n

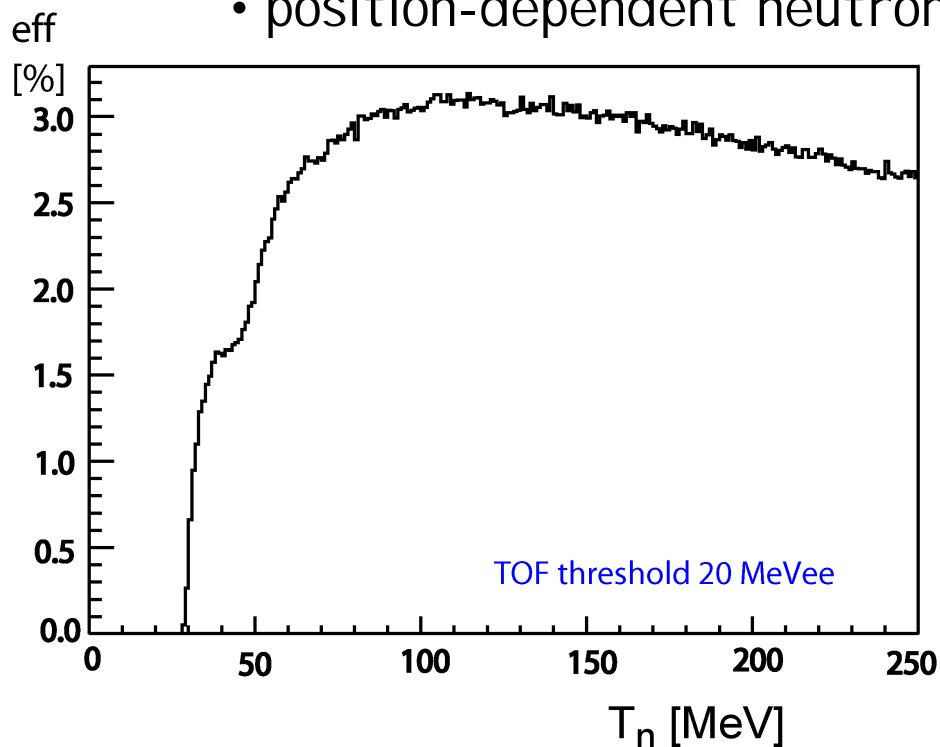
after accidental subtraction



integrate strength  
in  $-2 < E_m < 18$  MeV  
and plot versus  
other variables

# determination of experimental differential cross section

- luminosity
- dead-time effects
  - by means of (MC) simulations:
    - detection volume in phase space
    - hadronic interactions of proton
    - position-dependent neutron detection efficiency



effect of  
neutron threshold  
10→20 MeVee:  
Signal: factor 2 ↓↓  
Noise: factor 4 ↓↓

# calculations of ${}^3\text{He}(e, e'NN)$ cross section

continuum Faddeev calculations of the Bochum group  
(J. Golak et al., Phys. Rev. **C51** (1995) 1638)

Employ realistic NN-interactions like Bonn-B, CD-Bonn, V18, Nijmegen93,.

- parameter-free model
- calculate wave functions of  ${}^2\text{H}$  and  ${}^3\text{He}$ .
- include NN-rescattering up to all orders
- one-body current operator
- MECs included via a  $\pi$  and  $\rho$  exchange current operator (Schiavilla&Riska)
- $\Delta$ -current in, static,

low-energy approximation 
$$j_{p\Delta} \cong \frac{f_p^2}{(m_\Delta - m_N)} G_M^V(q)[\dots] \times q$$

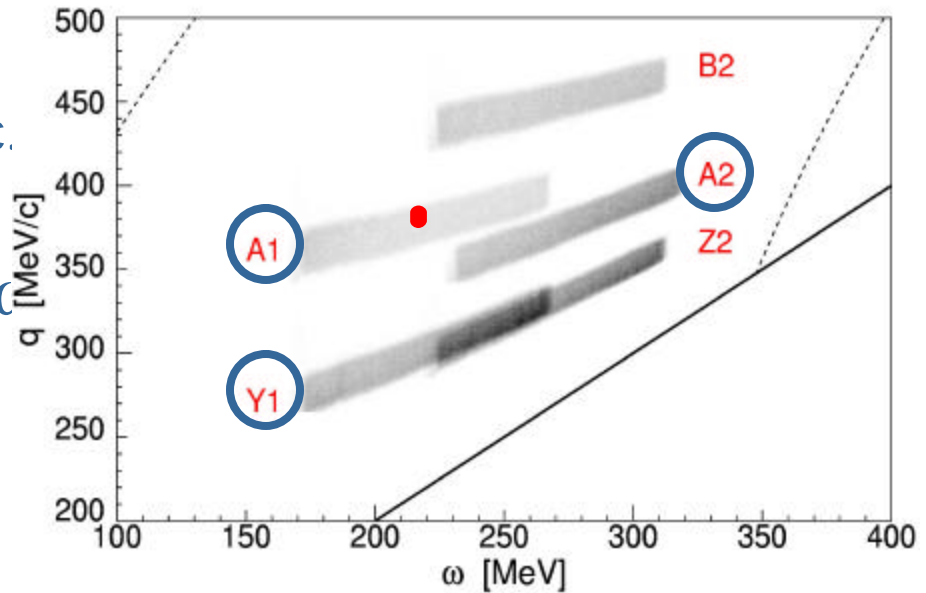
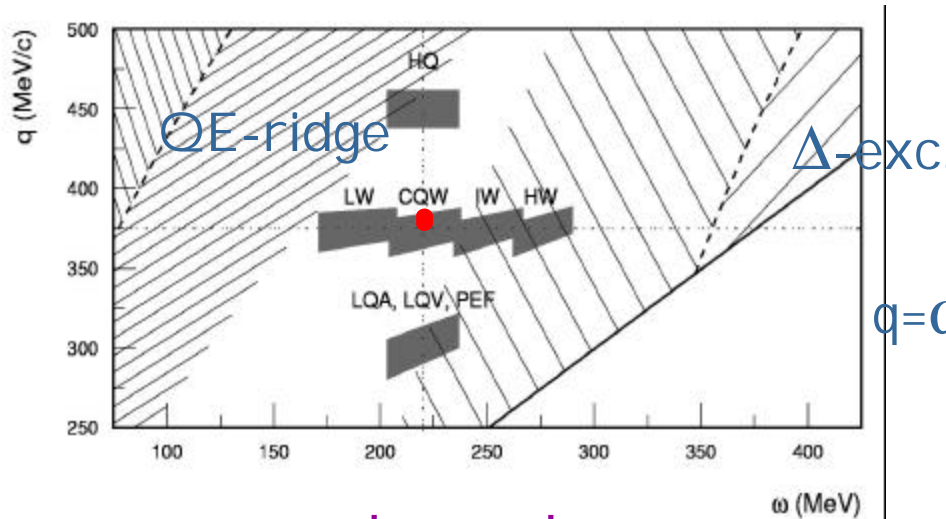
Here: Bonn-B potential,  $j \leq 3$ ,  $J \leq 15/2$

$\sigma_{\text{theo}}$  evaluated over  $2.5 \times 10^6$  grid points /kinematics  
in 5-D space of detected hadrons at the central value of  $(q, \omega)$

# Kinematic coverage

$^3\text{He}(e, e'pp)$  @ AmPS

$^3\text{He}(e, e'pn)$  @ MAMI



- $p_m$ -dependence
- $q$ -dependence
- $\omega$ -dependence

$d^8\sigma$  in  $\text{zm}^2/\text{MeV}^2/\text{sr}^3$   
 (1  $\text{zm}^2 = 10^{-14}$  barn)



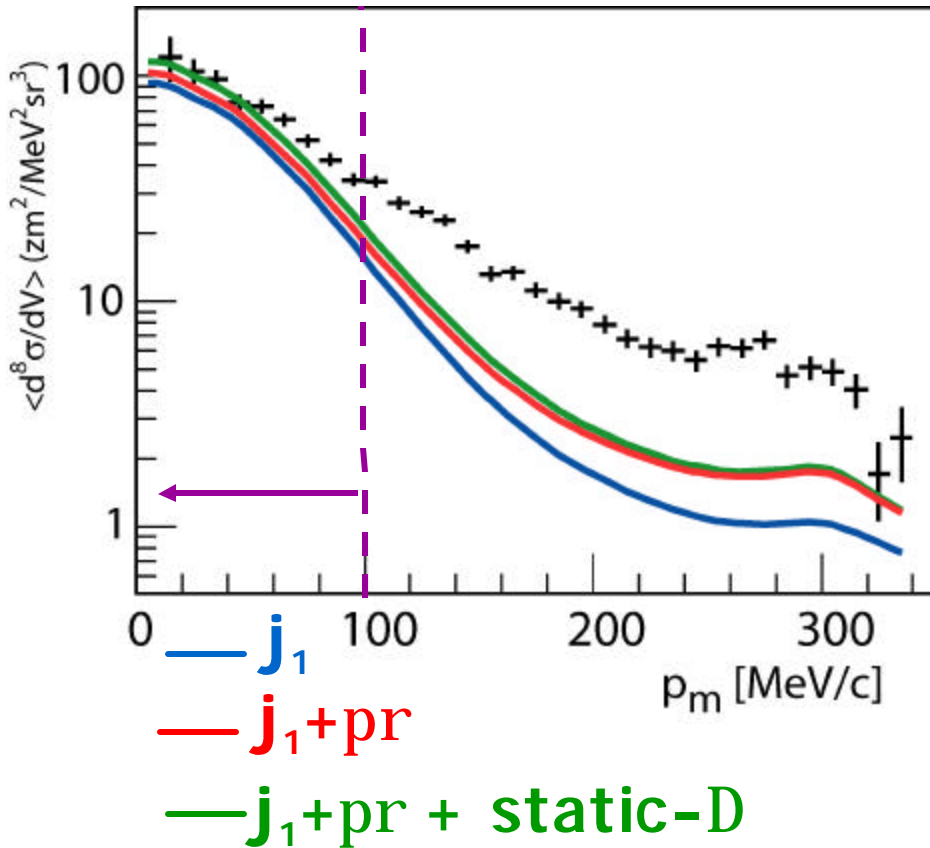
$^3\text{He}(e, e'pp)$

$p_m$  dependence

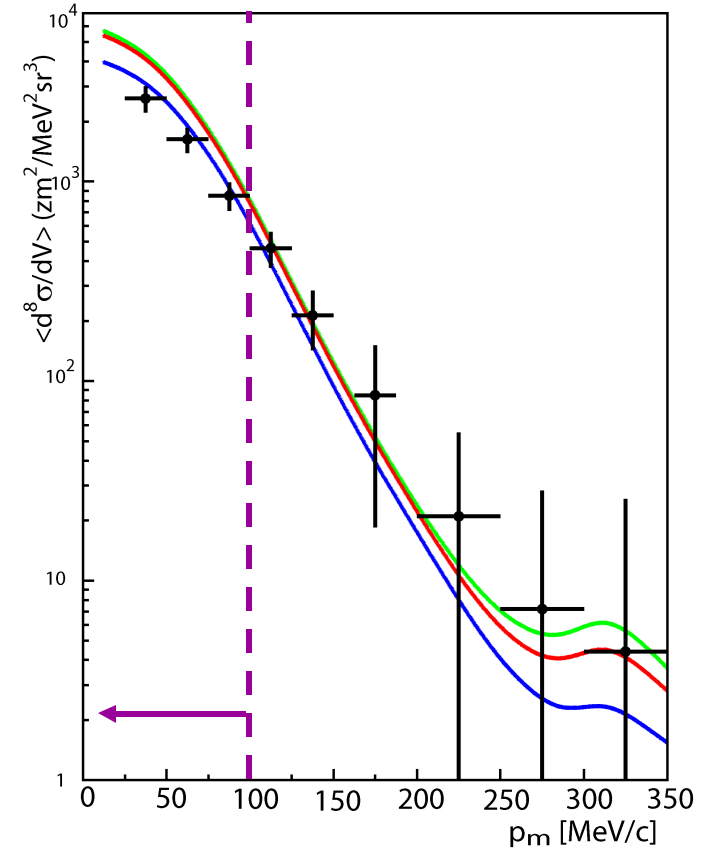
$q=300 \text{ MeV}/c \quad \omega=220 \text{ MeV}$

PRELIMINARY

$^3\text{He}(e, e'pn)$



$\frac{\text{data}}{\text{theory}} \approx 1.3$  for  $0 < p_m < 80 \text{ MeV}/c$   
 $\approx 3$  for  $p_m > 200 \text{ MeV}/c$



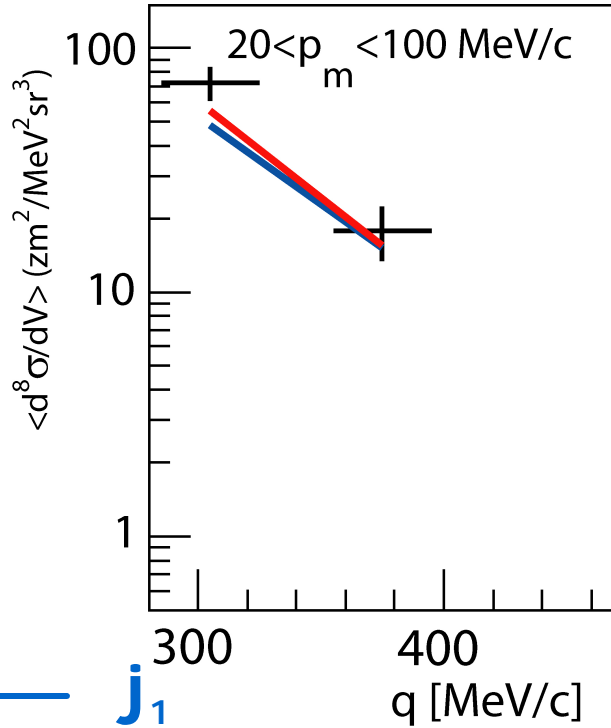
$\frac{\text{data}}{\text{theory}} \approx 0.6$  for  $p_m < 80 \text{ MeV}/c$

# q-dependence

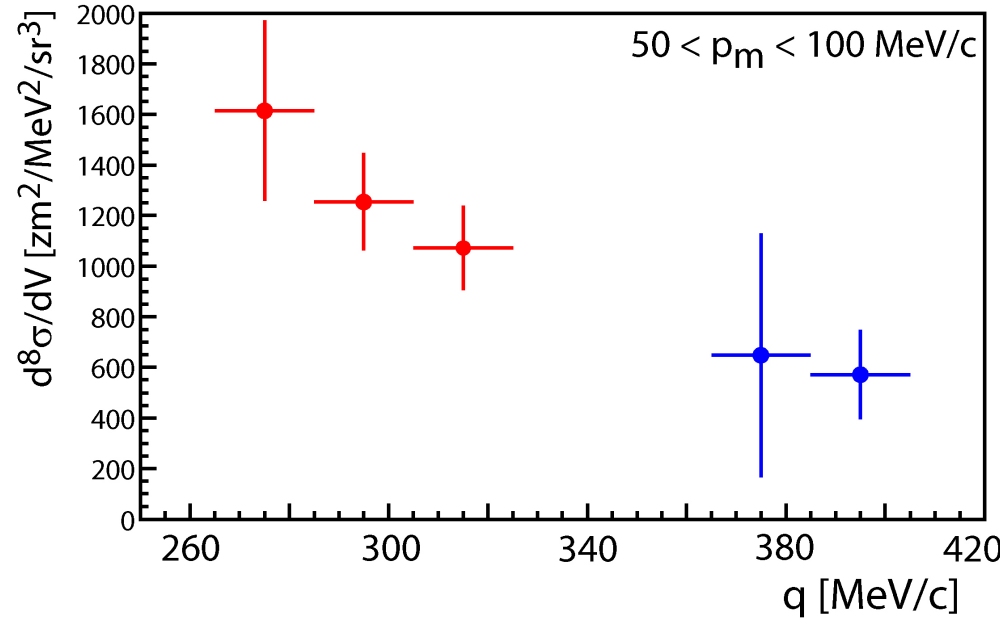
PRELIMINARY

${}^3\text{He}(e, e'pp)$

${}^3\text{He}(e, e'pn)$



$\omega = 220 \text{ MeV}$



—  $j_1$ 
—  $j_1 + pr$ 

$$\frac{s@(q=300)}{s@(q=375)} \gg 6$$

—  $s@(q=300)$ 
—  $s@(q=375)$ 

$$\frac{s@(q=300)}{s@(q=375)} \gg 2$$

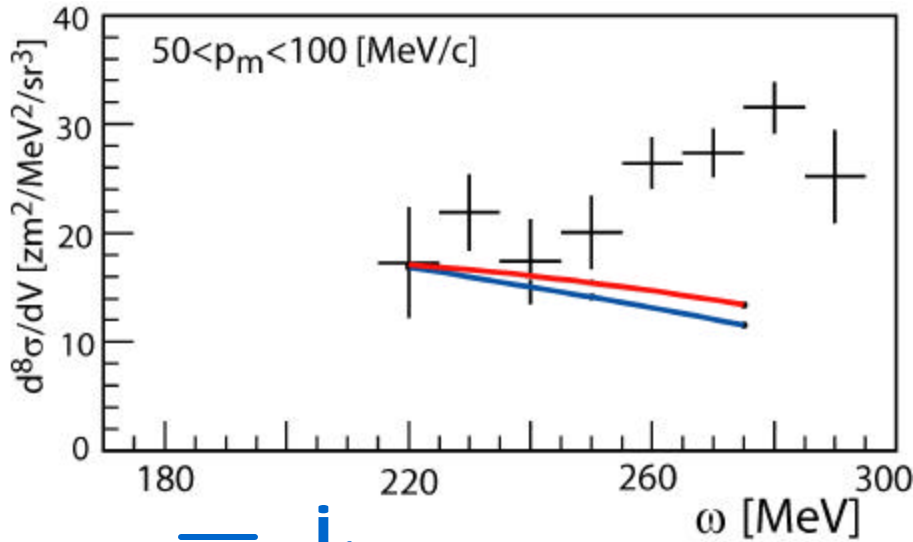
at low  $p_m$  the  ${}^3\text{He}(e, e'pp)$  reaction is dominated by direct two-proton emission induced by a one-body hadronic current.

$^3\text{He}(e, e'pp)$

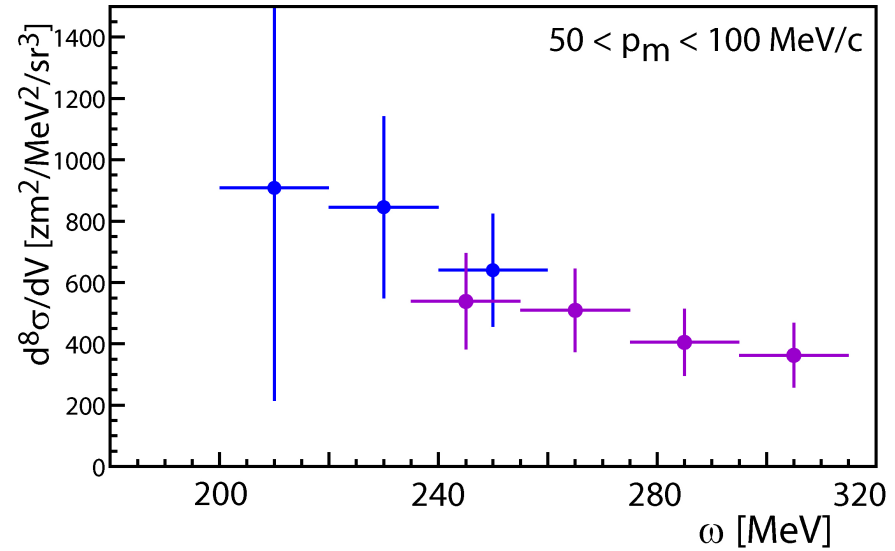
w dependence

$^3\text{He}(e, e'pn)$

$q=375 \text{ MeV}/c$



—  $j_1$   
—  $j_1 + pr$



$\frac{\text{data}}{\text{theory}} = 1.0$  for  $\omega \sim 220 \text{ MeV}$   
 $2.3$  for  $\omega \sim 280 \text{ MeV}$

$\frac{\text{data @ } (w=220)}{\text{data @ } (w=300)} \gg 2 (\pm 1)$

# Conclusions of ${}^3\text{He}(e, e'pp)$ and ${}^3\text{He}(e, e'pn)$

- at low  $p_m$  the  ${}^3\text{He}(e, e'pp)$  reaction seems well suited to study initial-state correlations.
- at  $200 < p_m < 350$  MeV/c continuum Faddeev calculations underestimate the  ${}^3\text{He}(e, e'pp)$  data by up to a factor 3.
- this discrepancy increases for increasing  $\omega$  (170  $\rightarrow$  290 MeV).
- ${}^3\text{He}(e, e'pn)$  is harder but feasible: at low  $p_m$  the measured cross sections are about a factor 2 smaller than predicted.
- q-dependence of  ${}^3\text{He}(e, e'pn)$  differs from that of pp-knockout.
- measured  ${}^3\text{He}(e, e'pn)$  cross section is decreasing between  $\omega=210$  and 320 MeV.
- finalize the  ${}^3\text{He}(e, e'pn)$  analysis.
- the role of the  $\Delta$ -resonance seriously needs theoretical attention. (i.e. Hannover and Bochum group).
- $(e, e'pN)$  data need to be measured with good statistics over a wide kinematic range and then sorted in narrow bins in many variables ( $q, \omega, p_m, \gamma_1, p_{ij}, \dots$ )